

初期宇宙における軽いダークフォトン ダークマターの生成

北嶋直弥



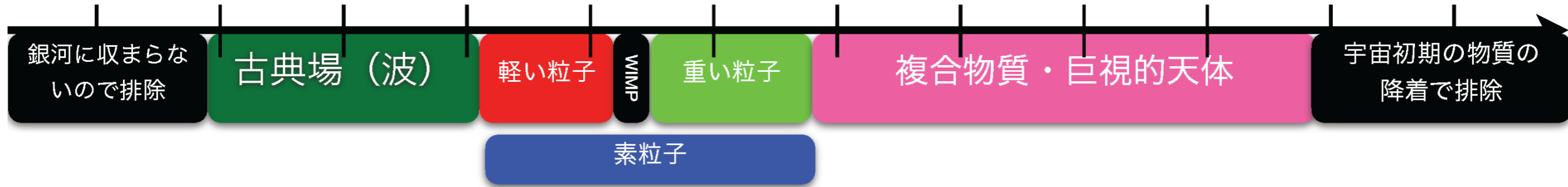
NK, Kazunori Nakayama (Tohoku U.), 2212.13573, 2303.04287
NK, Fuminobu Takahashi (Tohoku U.), 2303.05492

素粒子現象論研究会 2022, 3/16-18, 大阪公立大学

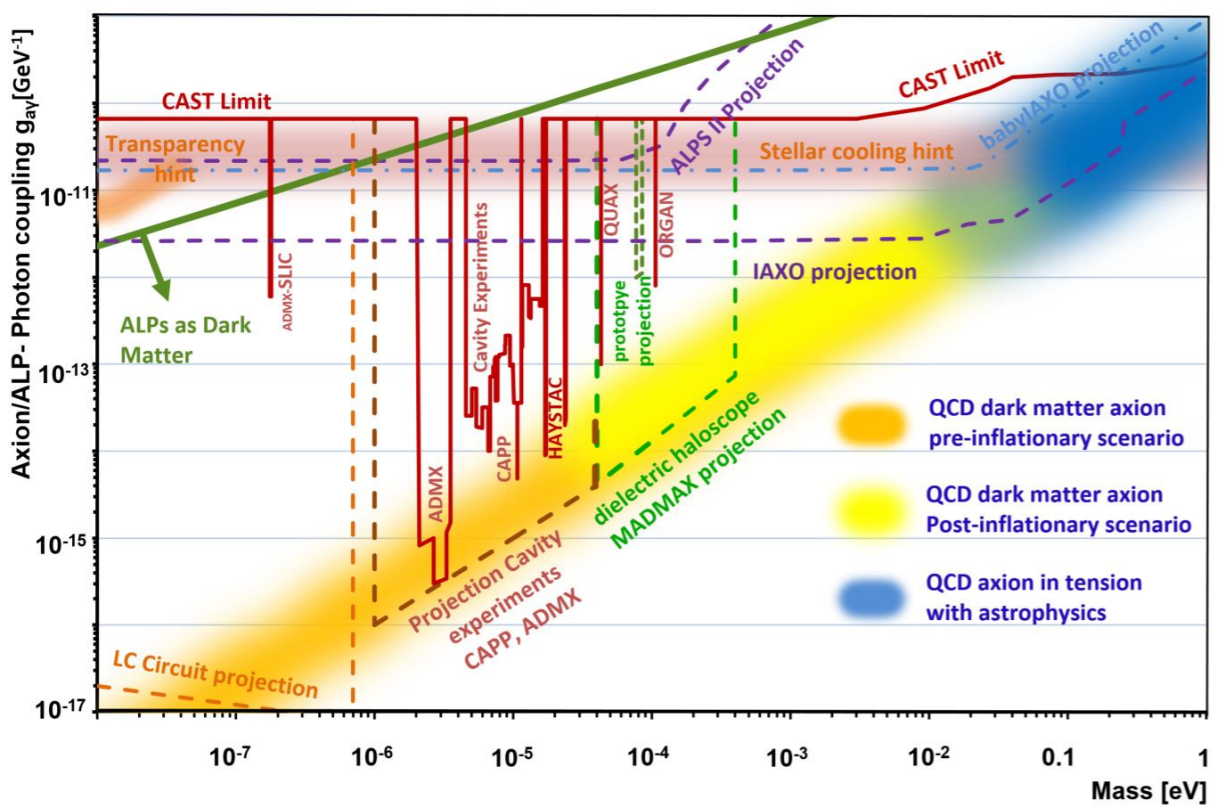
ニュートリノ 電子陽子 細菌 蚊 人間 富士山 地球 太陽

ダークマターの質量 [GeV/c²]

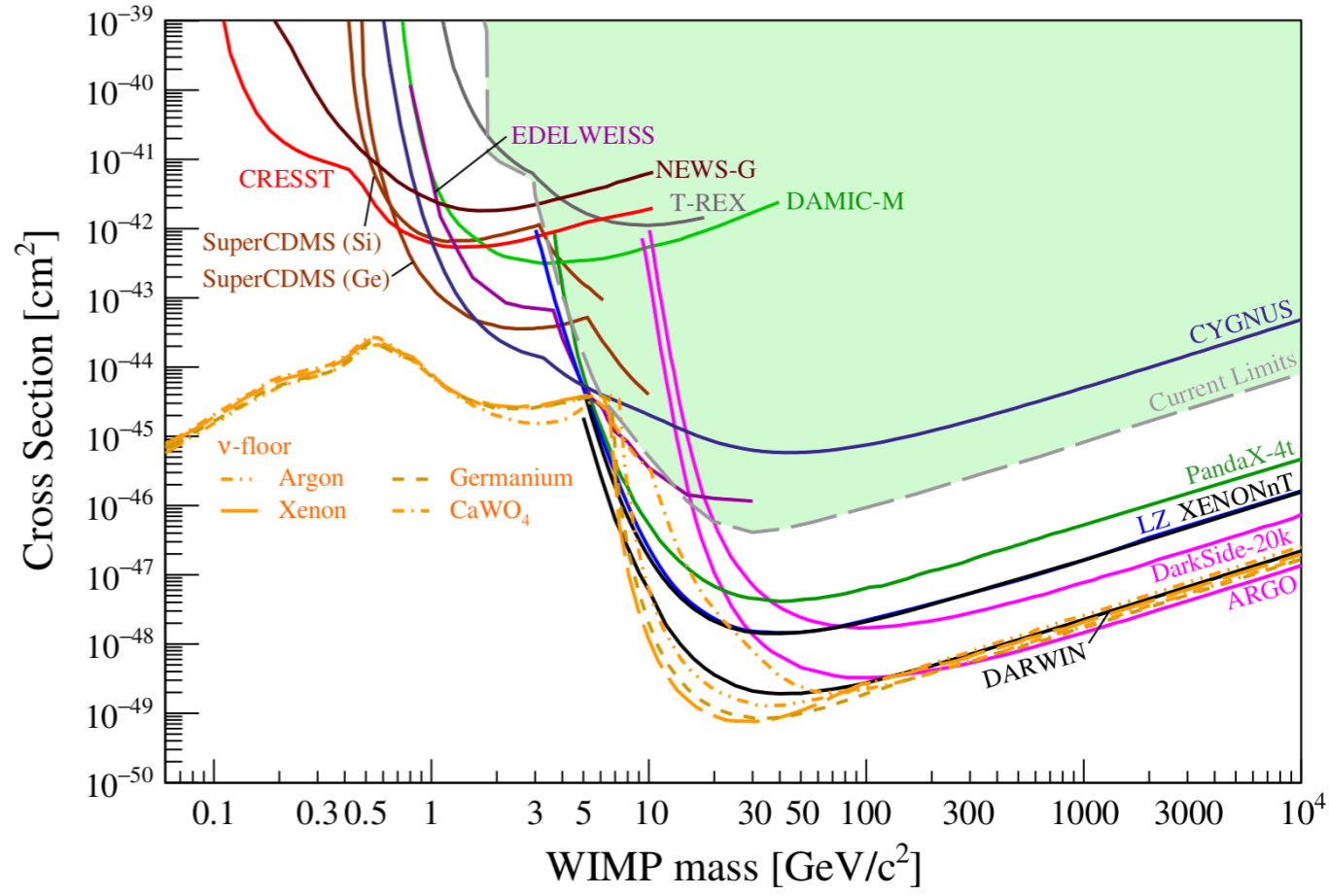
10⁻⁴⁰ 10⁻³⁰ 10⁻²⁰ 10⁻¹⁰ 10⁰ 10¹⁰ 10²⁰ 10³⁰ 10⁴⁰ 10⁵⁰ 10⁶⁰ 10⁷⁰



Axion search

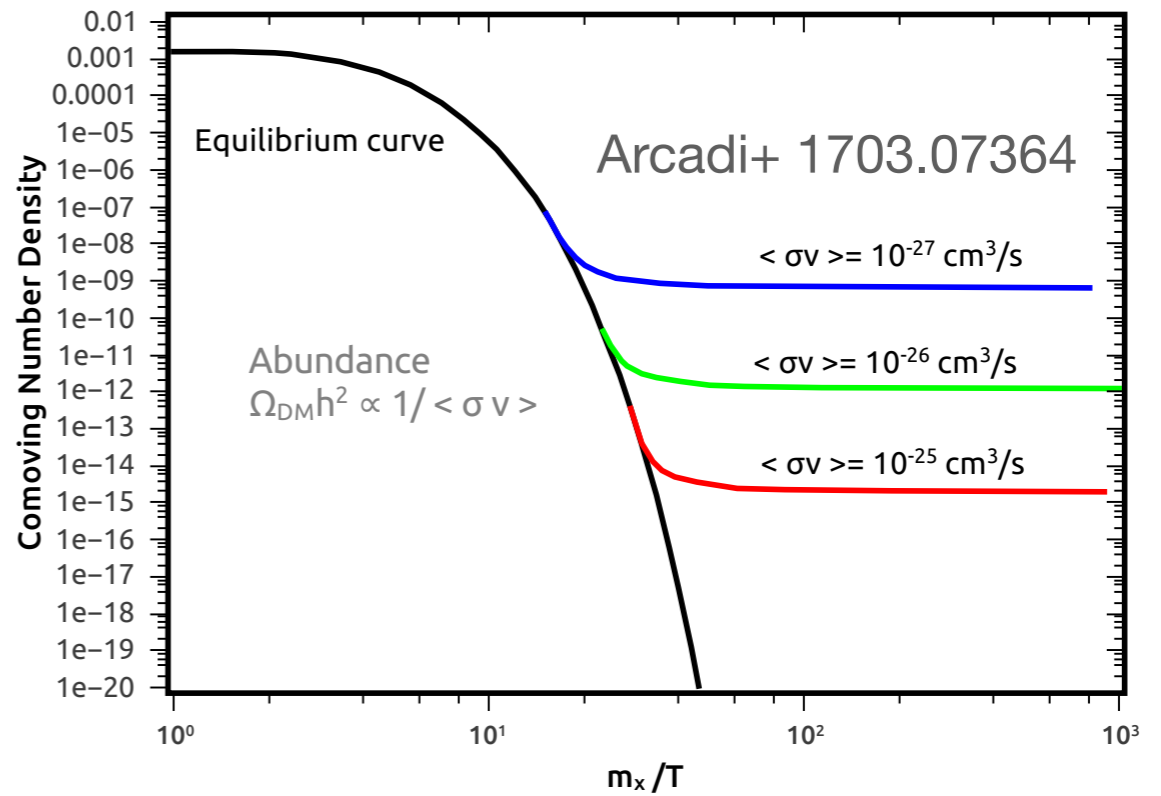
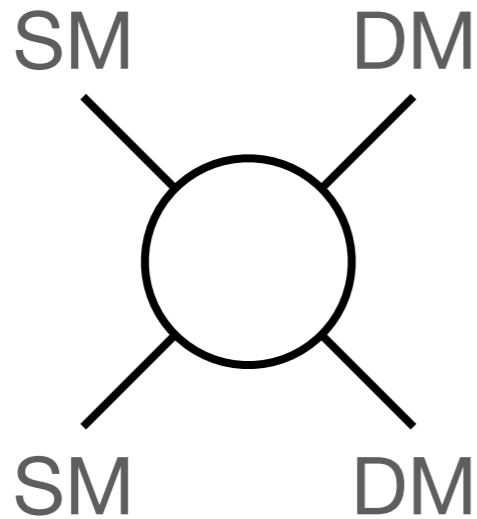


WIMP search



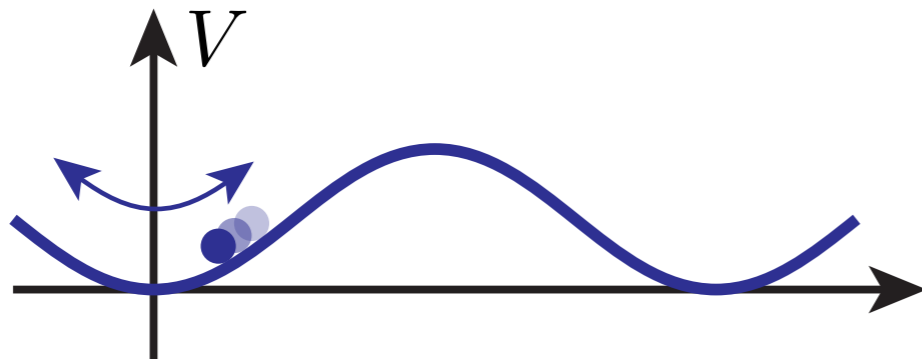
Dark matter production

WIMP DM : thermal freeze-out

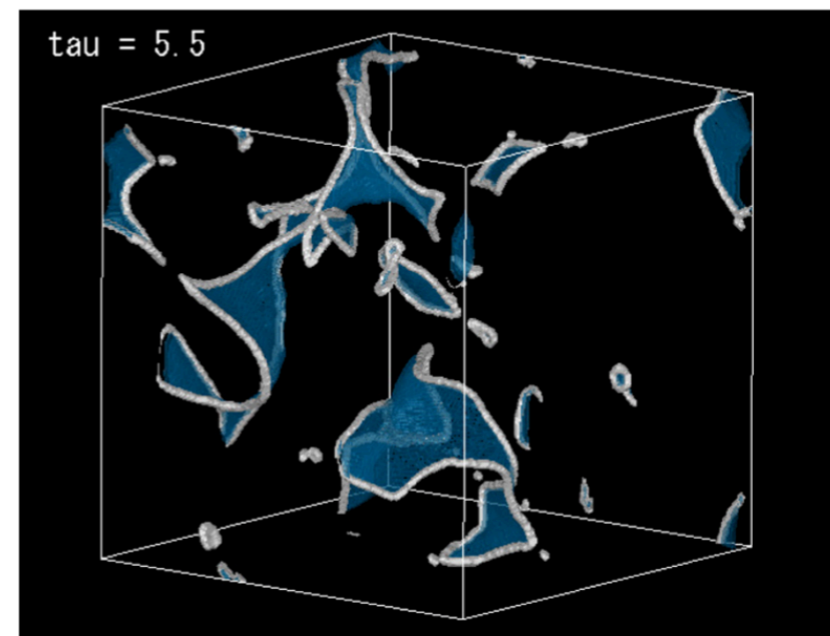


Axion DM : non-thermal production

- misalignment production



- topological defect



Hiramatsu et al 1202.5851

Dark photon DM production

- Gravitational particle production during inflation / reheating

Graham, Mardon, Rajendran (2016) / Ema, Nakayama, Tang (2019)

$$\Omega_{\gamma'} \simeq \Omega_{\text{DM}} \sqrt{\frac{m_{\gamma'}}{6 \mu\text{eV}}} \left(\frac{H_{\text{inf}}}{10^{14} \text{ GeV}} \right)^2 \rightarrow \text{lower limit on dark photon mass}$$

- Resonant production from scalar field

Axion : Agrawal, NK, Reece, Sekiguchi, Takahashi (2020)

Co, Pierce, Zhang, Zhao (2019), Bastro-Gil, Santiago, Ubaldi, Vega-Morales (2019)

Higgs : Harigaya, Narayan (2019)

- Misalignment production Nakayama (2019), Nakayama (2020), NK, Nakayama (2023)

- Production from cosmic strings Long, Wang (2019), NK, Nakayama (2022)

Dark photon DM production

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Resonant dark photon production from axion

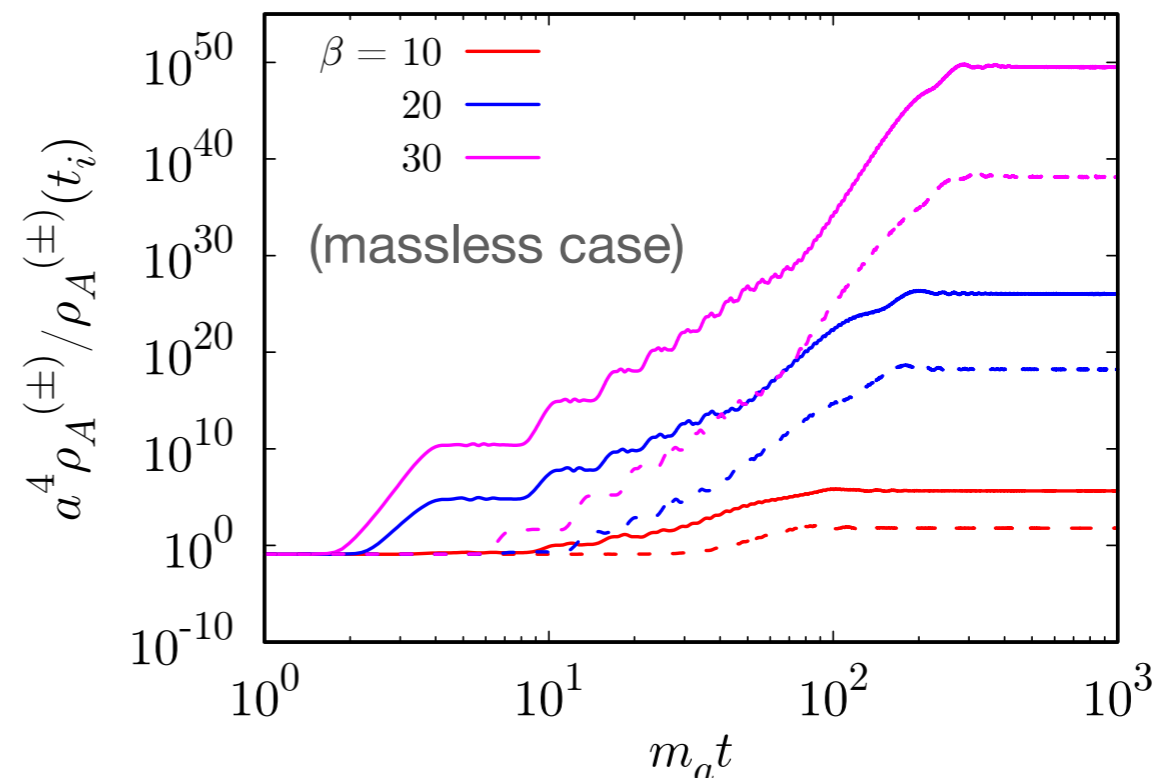
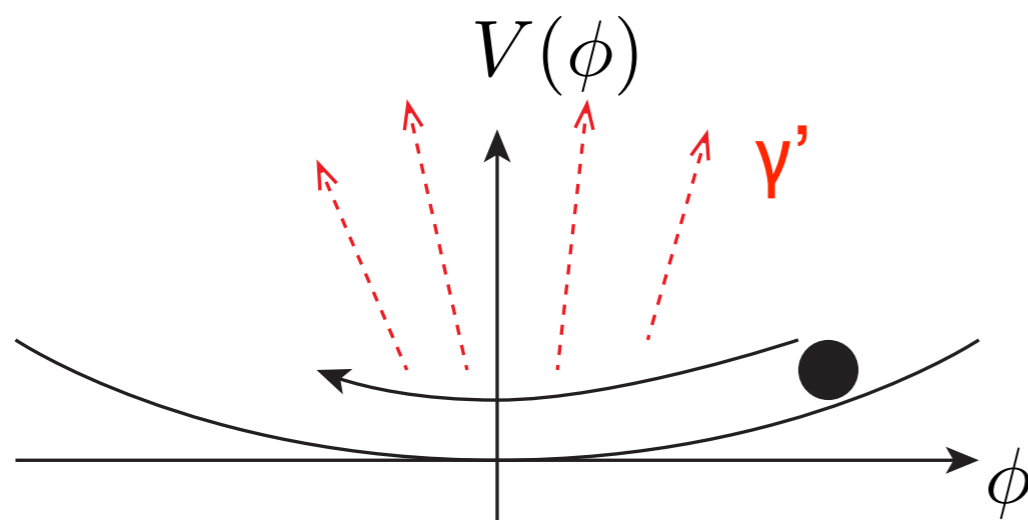
Agrawal, NK, Reece, Sekiguchi, Takahashi, 1810.07188

Co, Pierce, Zhang, Zhao, 1810.07196

Bastero-Gil, Santiago, Ubaldi, Vega-Morales, 1810.07208

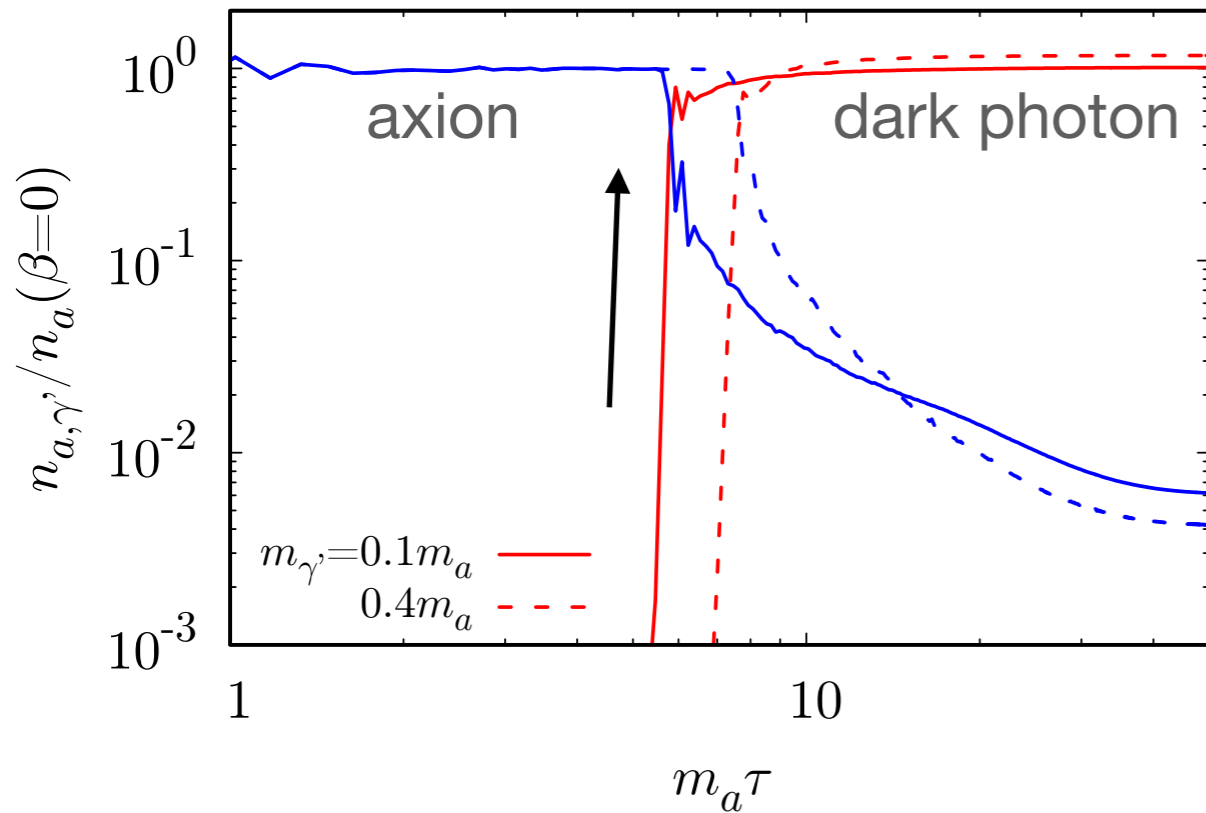
$$\mathcal{L} = \frac{1}{2} \partial^\mu \phi \partial_\mu \phi - V(\phi) - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_{\gamma'}^2 A_\mu A^\mu - \frac{\beta}{4f_a} \phi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

$$\longrightarrow \ddot{\mathbf{A}}_{\mathbf{k},\pm} + H \dot{\mathbf{A}}_{\mathbf{k},\pm} + \left(m_{\gamma'}^2 + \frac{k^2}{a^2} \mp \frac{k \beta \dot{\phi}}{a f_a} \right) \mathbf{A}_{\mathbf{k},\pm} = 0$$

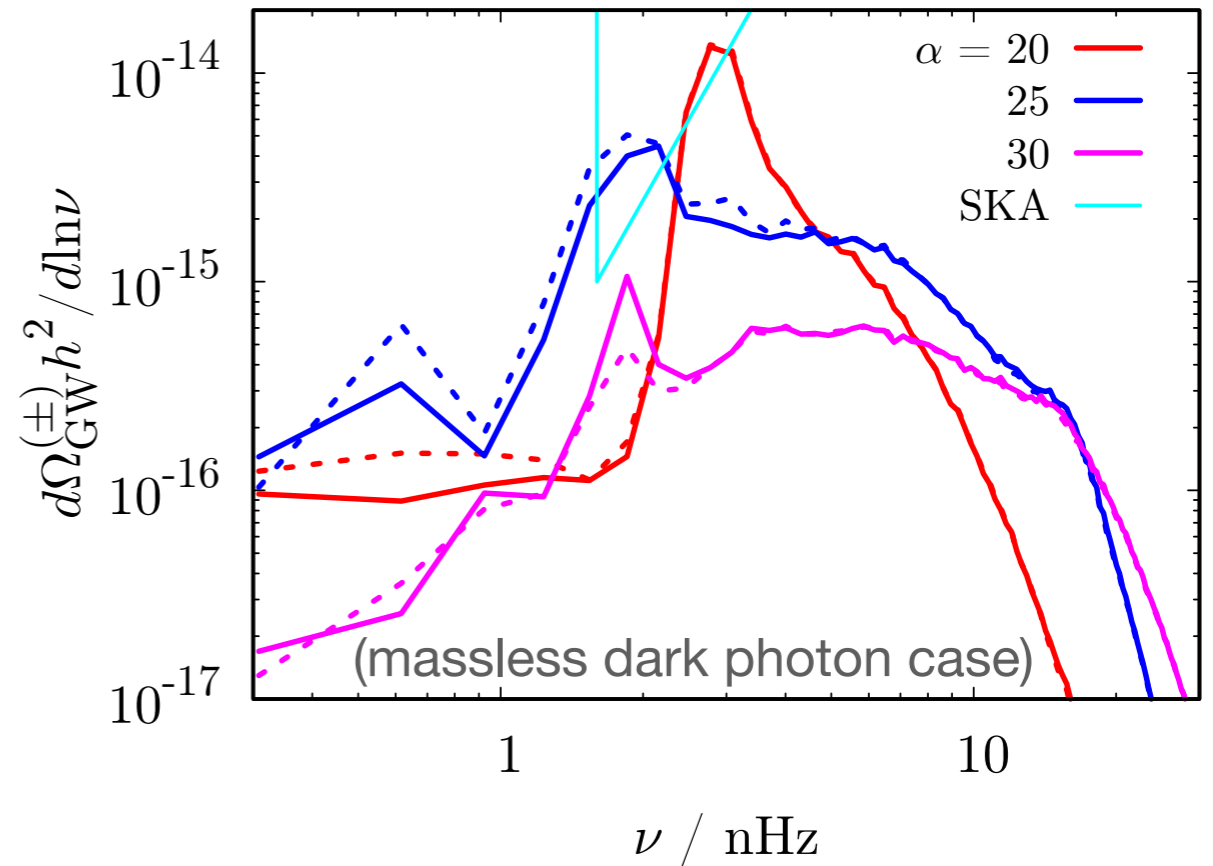


—> magnetogenesis Fujita+(2015), Kamada+(2019), Patel+(2020), ...

non-linear evolution



gravitational wave



- Axion abundance is suppressed & dark photon is dominant

Agrawal, NK, Reece, Sekiguchi, Takahashi, 1810.07188
(see also NK, T. Sekiguchi, F. Takahashi, 1711.06590)

- Produced dark photons can stabilize the dark Higgs $V(\Phi) \ni |\mathbf{A}|^2 |\Phi|^2$

—> secondary inflation, early dark energy

NK, Nakagawa, Takahashi, 2111.06696 Nakagawa, Takahashi, Yin, 2209.01107

- GW emission with circular polarization NK, Soda, Urakawa, 2010.10990

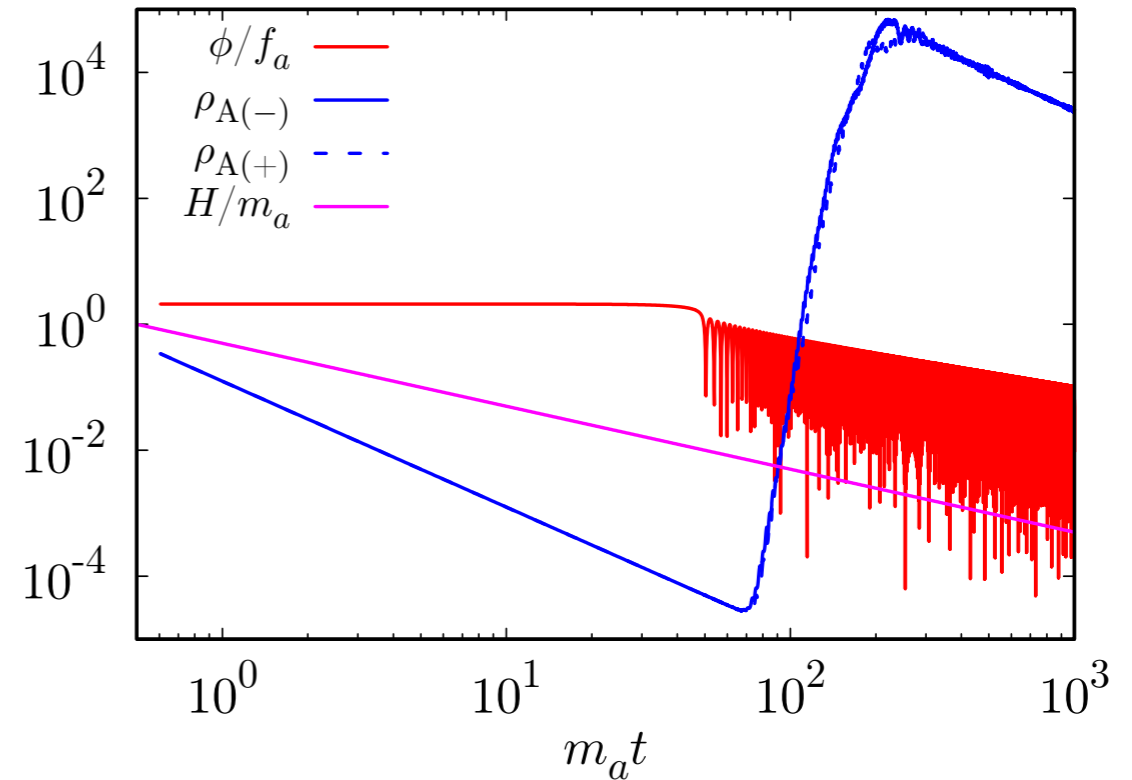
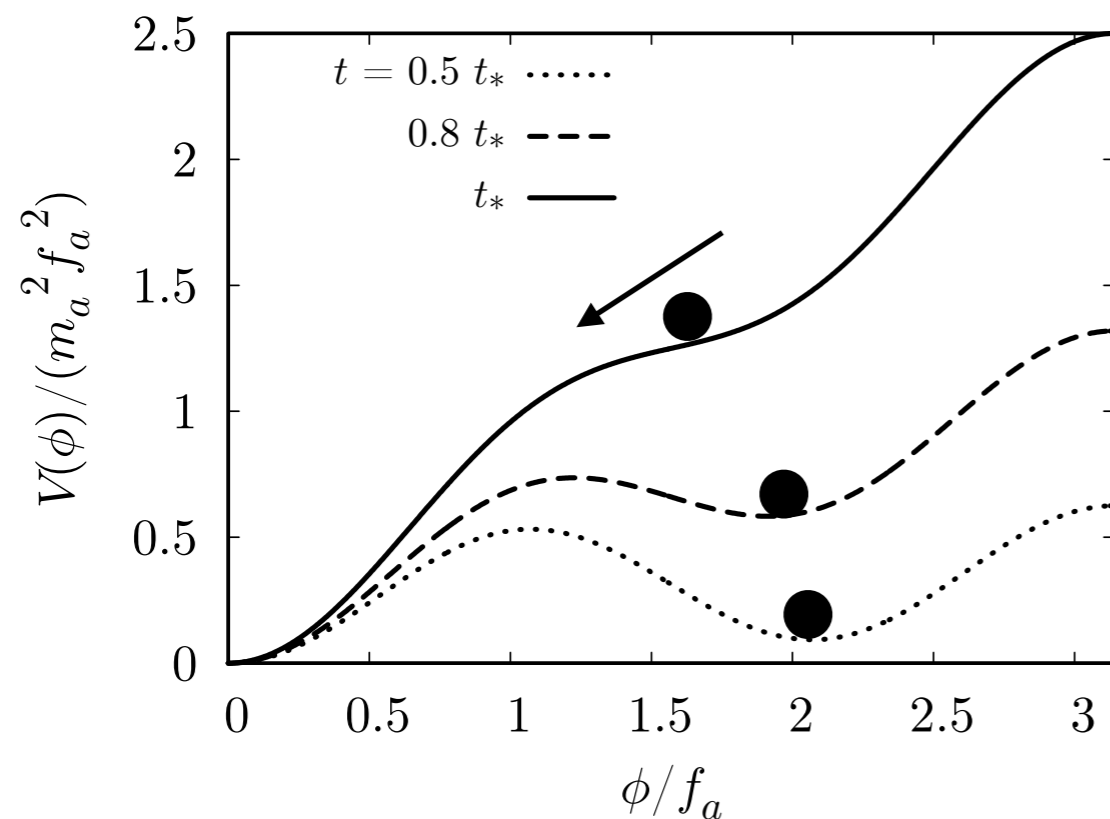
see also Machado+ (2019), Salehian+ (2020), Ratzinger+ (2020), Namba+ (2020)

Resonant dark photon production w/o large coupling

NK, Takahashi, 2303.05492

$$V(\phi) = m_a(t)^2 f_a^2 \left[1 - \cos \left(\frac{\phi}{f_a} \right) \right] + \Lambda_H^4 \left[1 - \cos \left(\frac{N_H \phi}{f_a} \right) \right]$$

$$m_a(t) = \begin{cases} m_{a0} (t/t_*)^{b/2} & \text{for } t < t_* \\ m_{a0} & \text{otherwise} \end{cases}$$

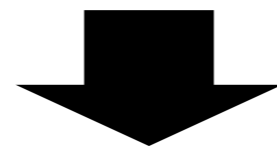


Application for QCD axion cosmology → next talk by Shota Nakagawa

Coherent vector DM production

Nakayama (2019), Nakayama (2020), NK, Nakayama (2023)

$$\mathcal{L} = -\frac{f^2(\phi)}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2} m_A^2 A_\mu A^\mu - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$$



$$f^2 \propto a^\alpha, \quad \bar{A}_i = f A_i / a, \quad R_A = \frac{\rho_A}{\rho_\phi}$$

$$\ddot{\phi} + 3H\dot{\phi} + \partial_\phi V \left(1 + \frac{\alpha R_A}{2\epsilon_V} \right) = 0 \quad \epsilon_V = \frac{M_P^2}{2} \left(\frac{\partial_\phi V}{V} \right)^2$$

(slow-roll parameter)

$$\ddot{\bar{A}}_i + 3H\dot{\bar{A}}_i + \left(\frac{m_A^2}{f^2} - \frac{(\alpha + 4)(\alpha - 2)}{4} H^2 + \frac{2 - \alpha}{2} \dot{H} \right) \bar{A}_i = 0$$

Statistical anisotropy $\mathcal{P}_\zeta(\mathbf{k}) = \mathcal{P}_\zeta^{(\text{iso})}(k)(1 + g_k \sin^2 \theta_k), \quad \hat{\mathbf{k}} \cdot \hat{\mathbf{A}} = \cos \theta_k$

$$\& \quad g_k \propto R_A$$

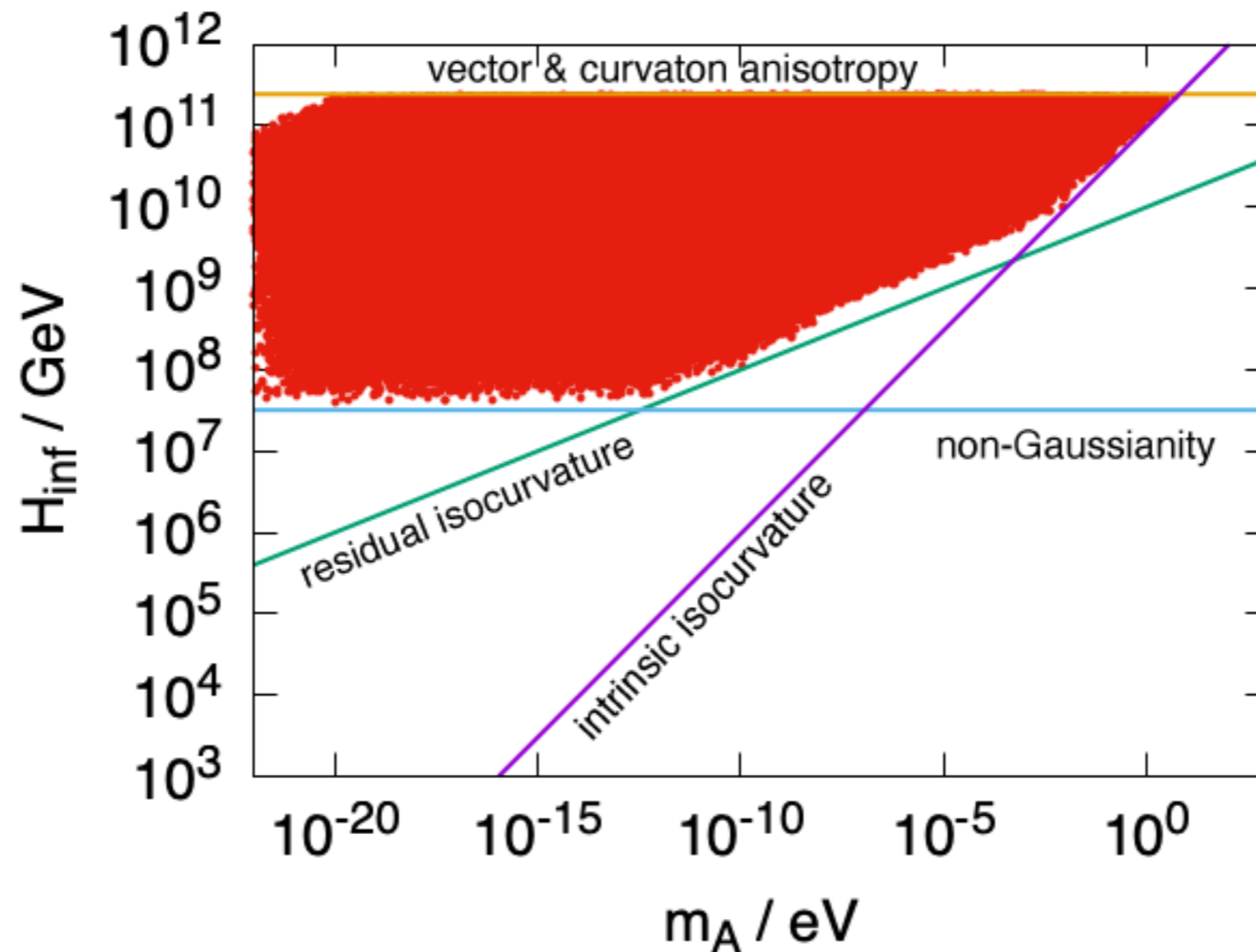
DM isocurvature perturbation $S = \frac{\delta \rho_A}{\bar{\rho}_A} \sim \frac{H_{\text{inf}}}{\pi \bar{A}_i} \propto R_A^{-1}$

CMB observation $\rightarrow g_k \lesssim 0.01, \quad S \lesssim 0.1\zeta$

“Viable” coherent vector DM scenario

NK, Nakayama, 2303.04287

curvaton scenario : introduction of an additional scalar field responsible for the curvature perturbation

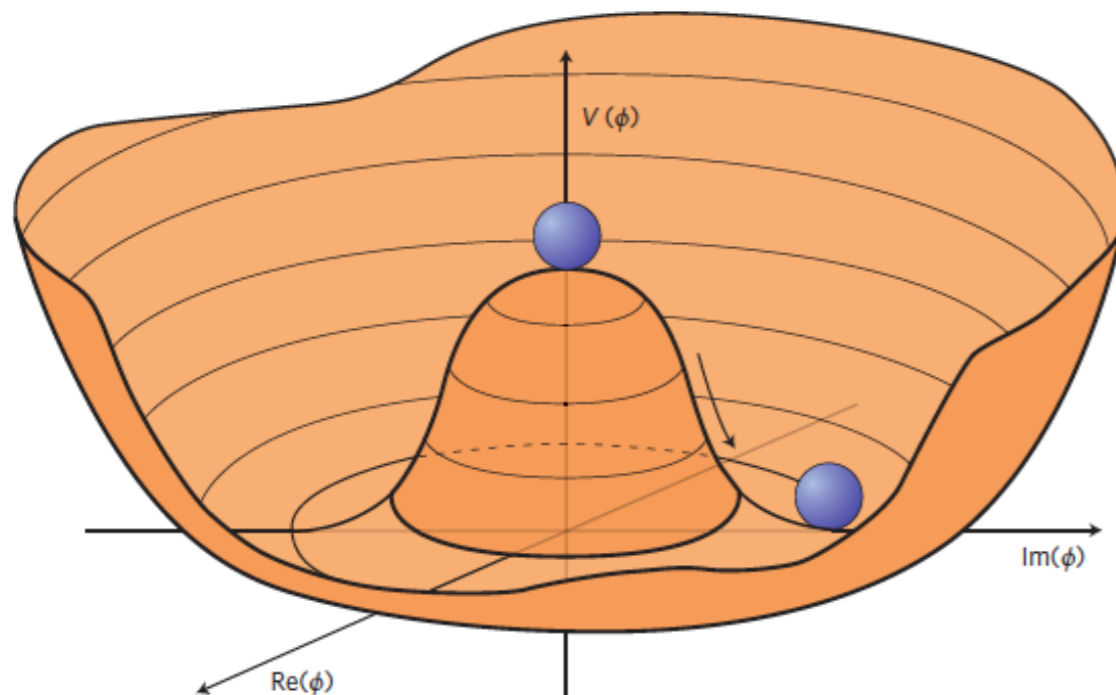


Dark photon DM from Abelian-Higgs cosmic strings

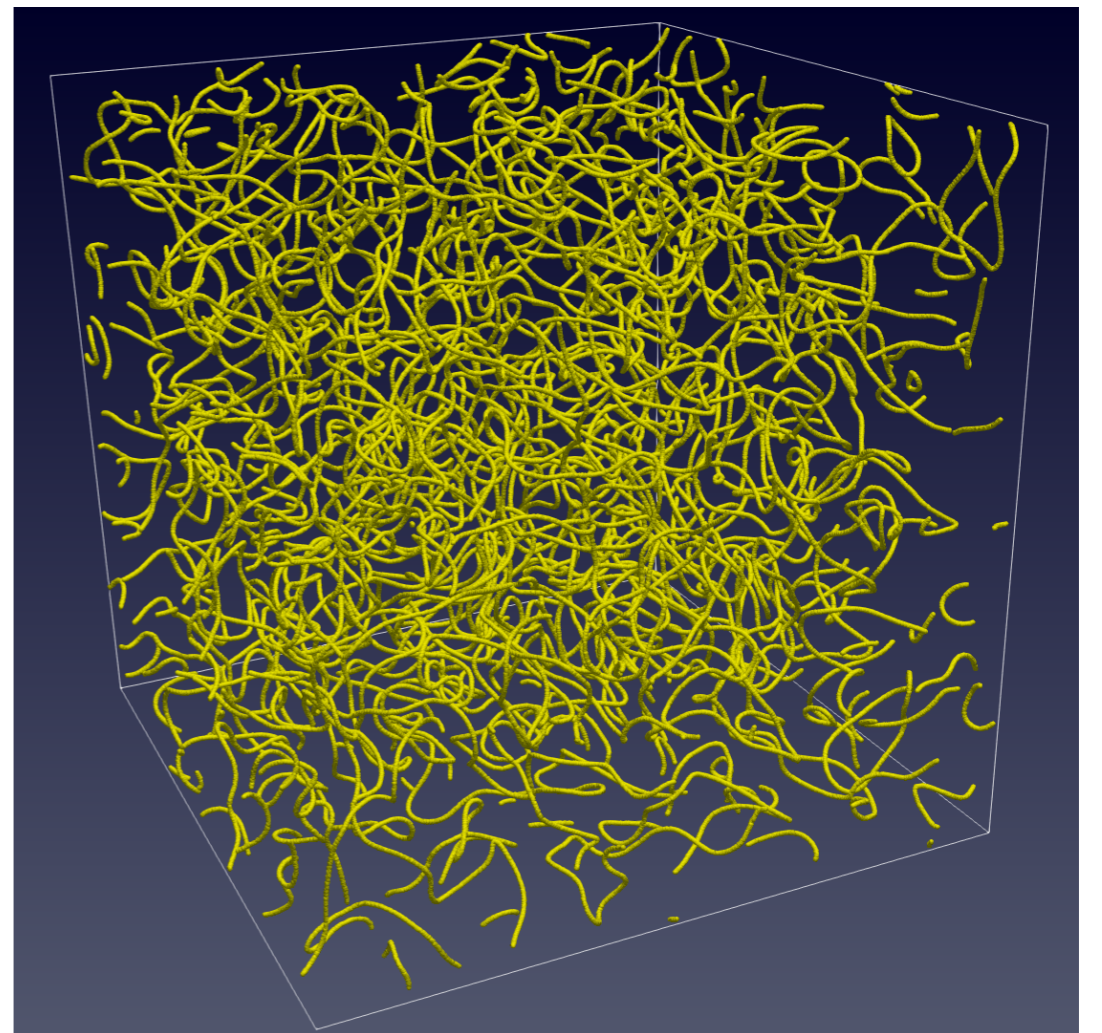
Long, Wang 1901.03312, NK, Nakayama 2212.13573

$$\mathcal{L} = (\mathcal{D}_\mu \Phi)^* \mathcal{D}^\mu \Phi - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - V(\Phi), \quad V(\Phi) = \frac{\lambda}{4} (|\Phi|^2 - v^2)^2$$
$$(\mathcal{D}_\mu = \partial_\mu - ieA_\mu, \quad F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu)$$

spontaneous U(1) symmetry breaking
—> formation of cosmic strings



Ellis, Gaillard, Nanopoulos 1504.07217



Scenario

- “Light” dark photons can be produced by cosmic strings

small gauge coupling

$e = 0$ limit corresponds to the axion emission (global string case)

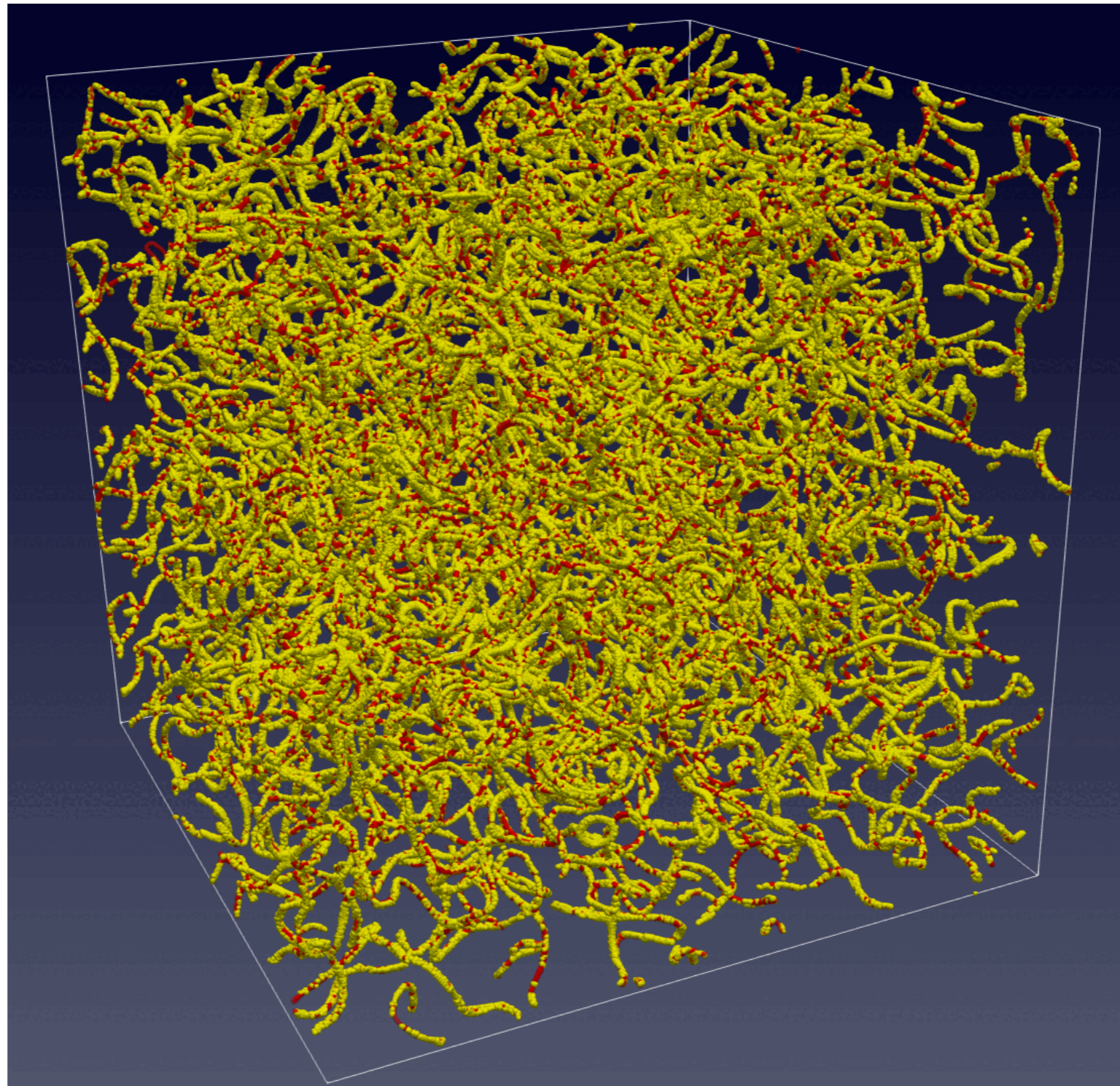
- Dark photon production becomes inefficient for $\ell_{\text{loop}} \gtrsim m_A^{-1}$
(i.e. loop oscillation frequency becomes smaller than the mass) $\rightarrow H \lesssim m_A$

- After that, string evolves like “local” string

(network loses the energy only through the GW emission)

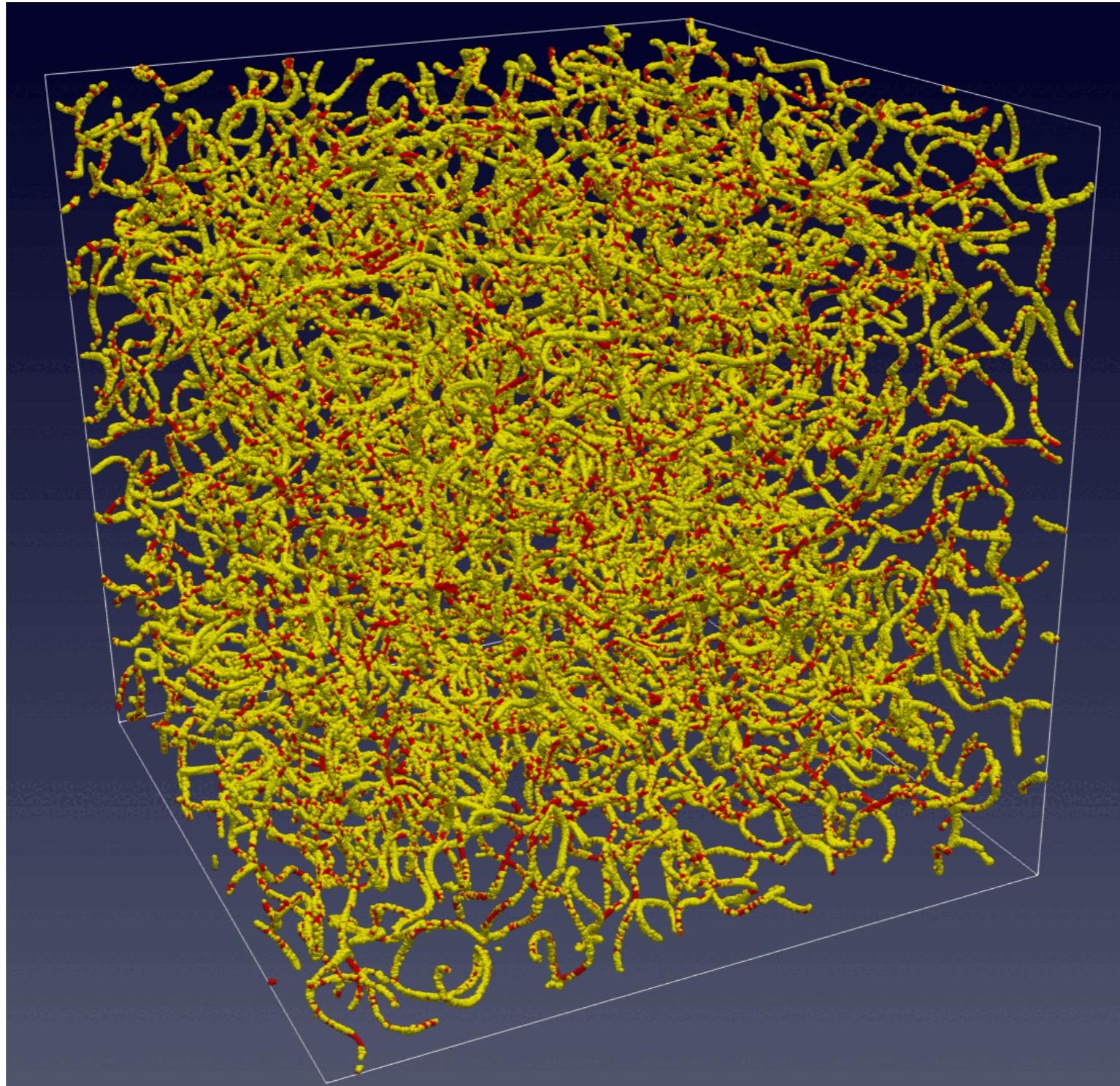
Loop production & decay

Type-II string with $e=0.01$ and $\lambda=2$

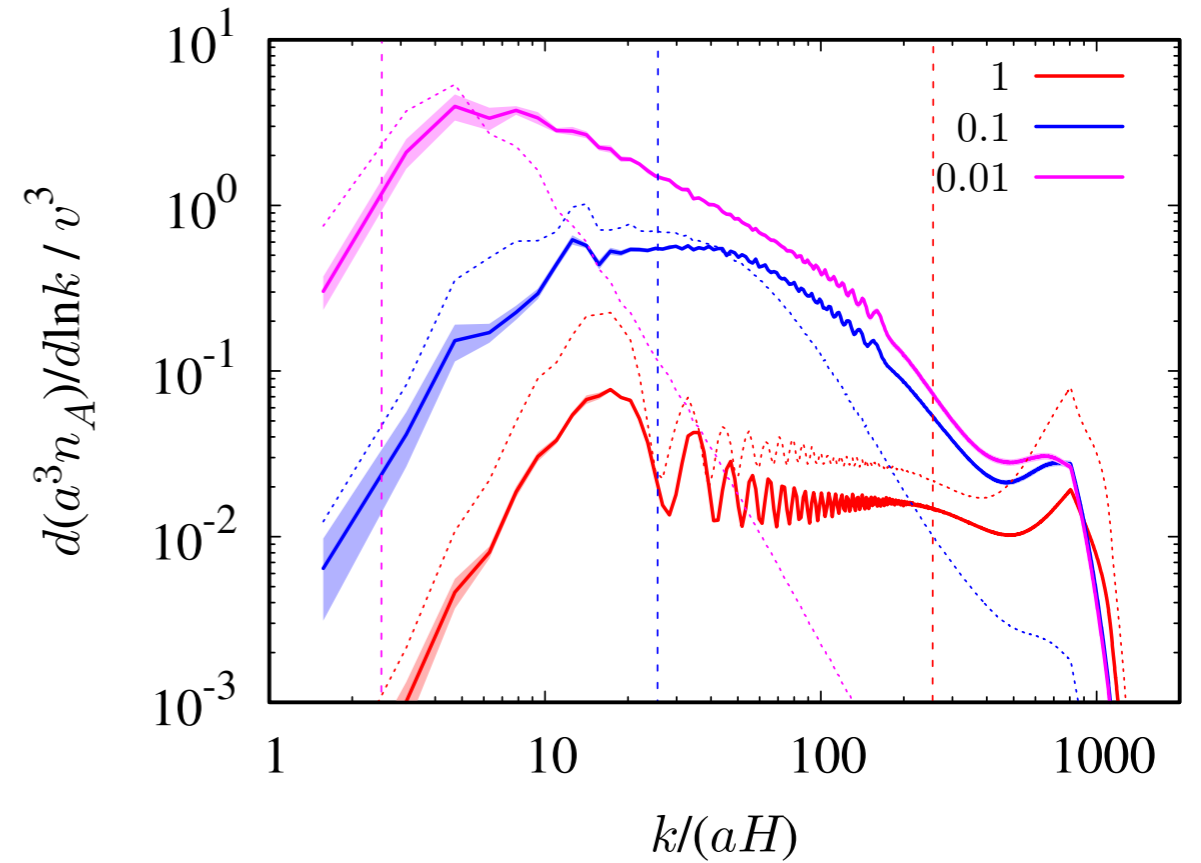
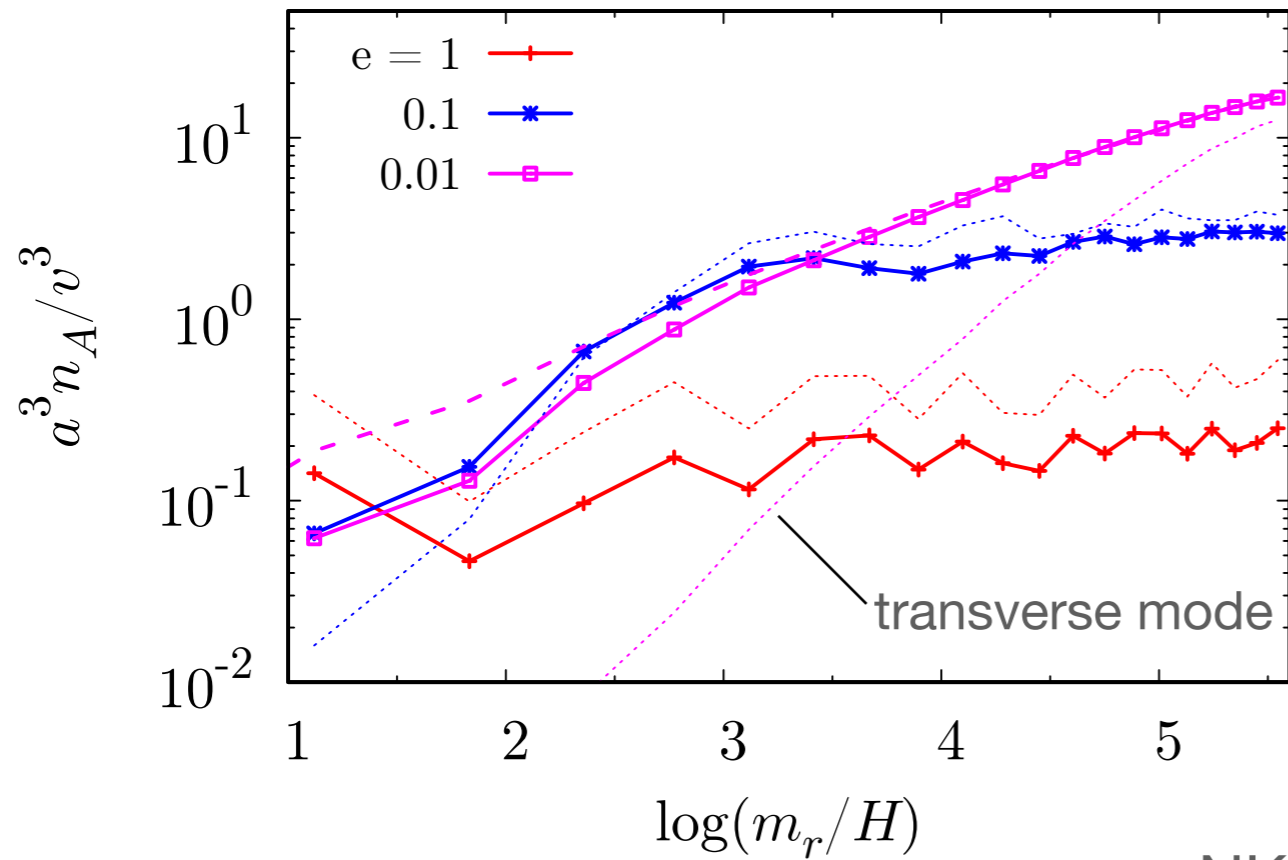


Loop production & decay

Type-I string with $e=1$ and $\lambda=2$



Dark photon DM abundance & spectrum



NK, Nakayama 2212.13573

$$\Omega_A h^2 = \frac{m_A (n_{A,0}/s_0) h^2}{\rho_{\text{cr},0}/s_0} \simeq 0.091 \left(\frac{\xi}{12} \right) \left(\frac{m_A}{10^{-13} \text{ eV}} \right)^{1/2} \left(\frac{v}{10^{14} \text{ GeV}} \right)^2$$

$$\xi = 0.15 \log \left(\frac{m_r}{m_A} \right) \simeq 12 + 0.15 \log \left[\left(\frac{m_r}{10^{14} \text{ GeV}} \right) \left(\frac{10^{-13} \text{ eV}}{m_A} \right) \right]$$

GW emission from cosmic strings



Credit: Daniel Dominguez/CERN

Quadrupole formula for GW emission: $\dot{E}_{\text{GW}} \sim G(\ddot{D})^2$

quadrupole moment: $D \sim ML^2 \sim \mu L^3$, $\ddot{D} \sim \omega^3 D \sim L^{-3} D$

L : typical loop size \sim (typical oscillation frequency)⁻¹

GW emission rate: $\dot{E}_{\text{GW}} \sim G\mu^2 \equiv \Gamma_{\text{GW}} G\mu^2$

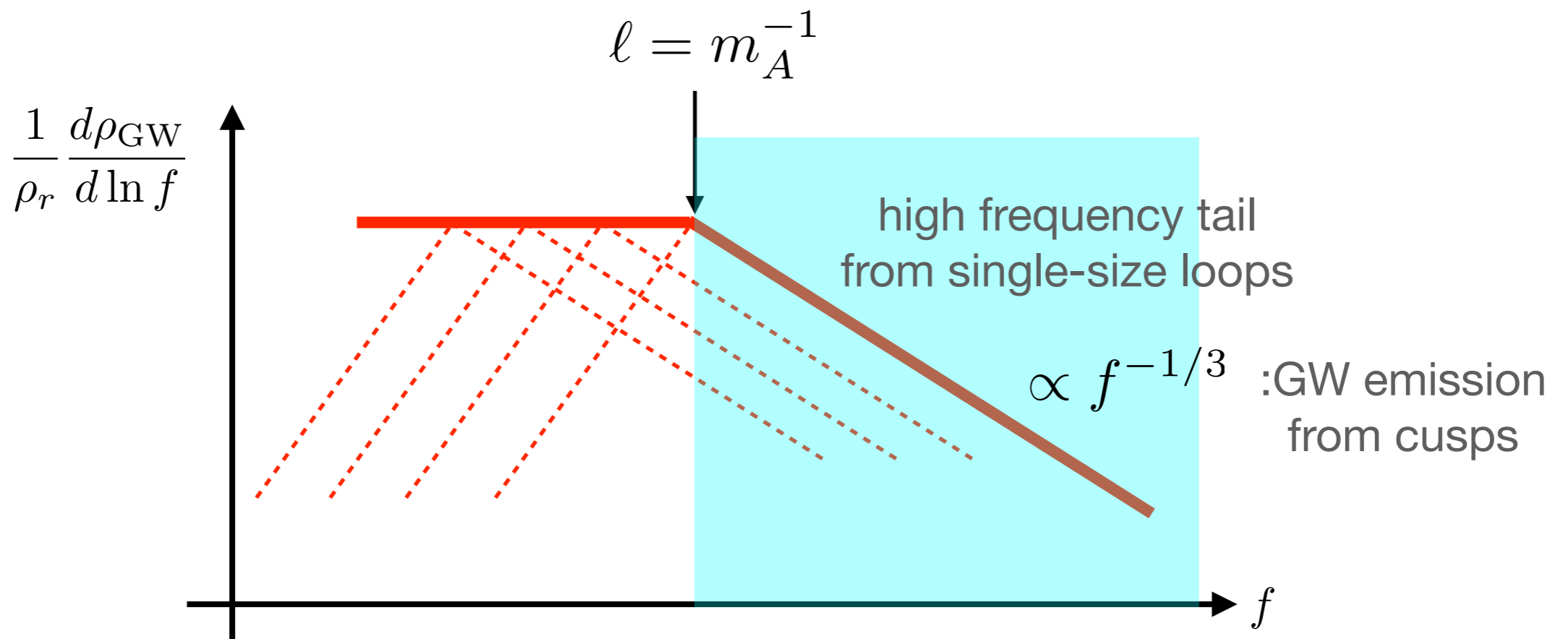
$G\mu \sim (v/M_P)^2 \sim 10^{-7} (v/10^{15} \text{ GeV})^2$

Energy loss of loops = GW emission + vector boson emission

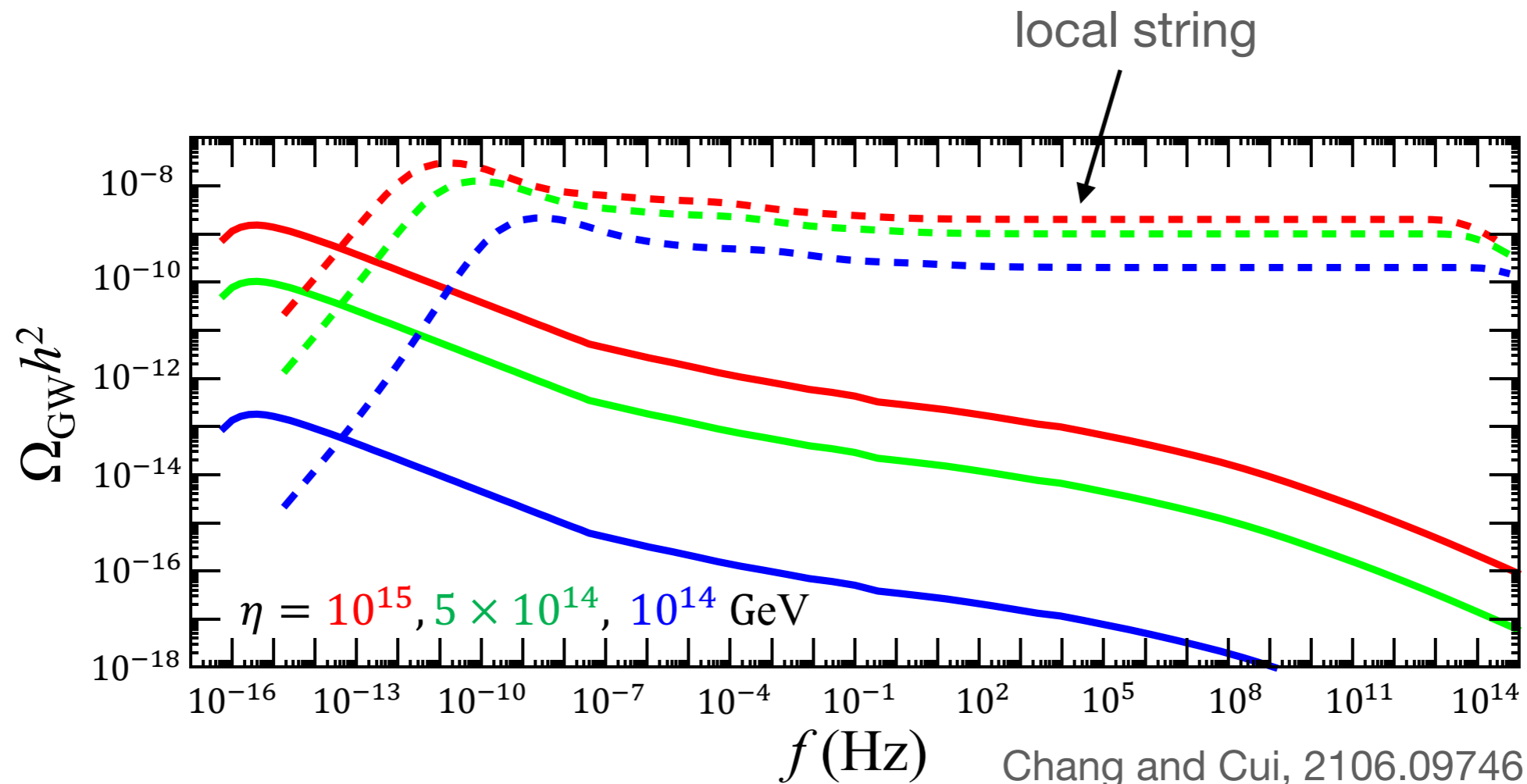
$$\frac{dE_\ell}{dt} = -\Gamma_{\text{GW}} G\mu^2 - \Gamma_{\text{vec}} v^2 \theta(1 - m_A \ell) \quad (\Gamma_{\text{GW}} \sim \Gamma_{\text{vec}} \sim 50)$$

Loops shorter than m_A^{-1} can emit dark photons

→ short lived & GW emission is suppressed



GW spectrum from local/global strings (conventional models)

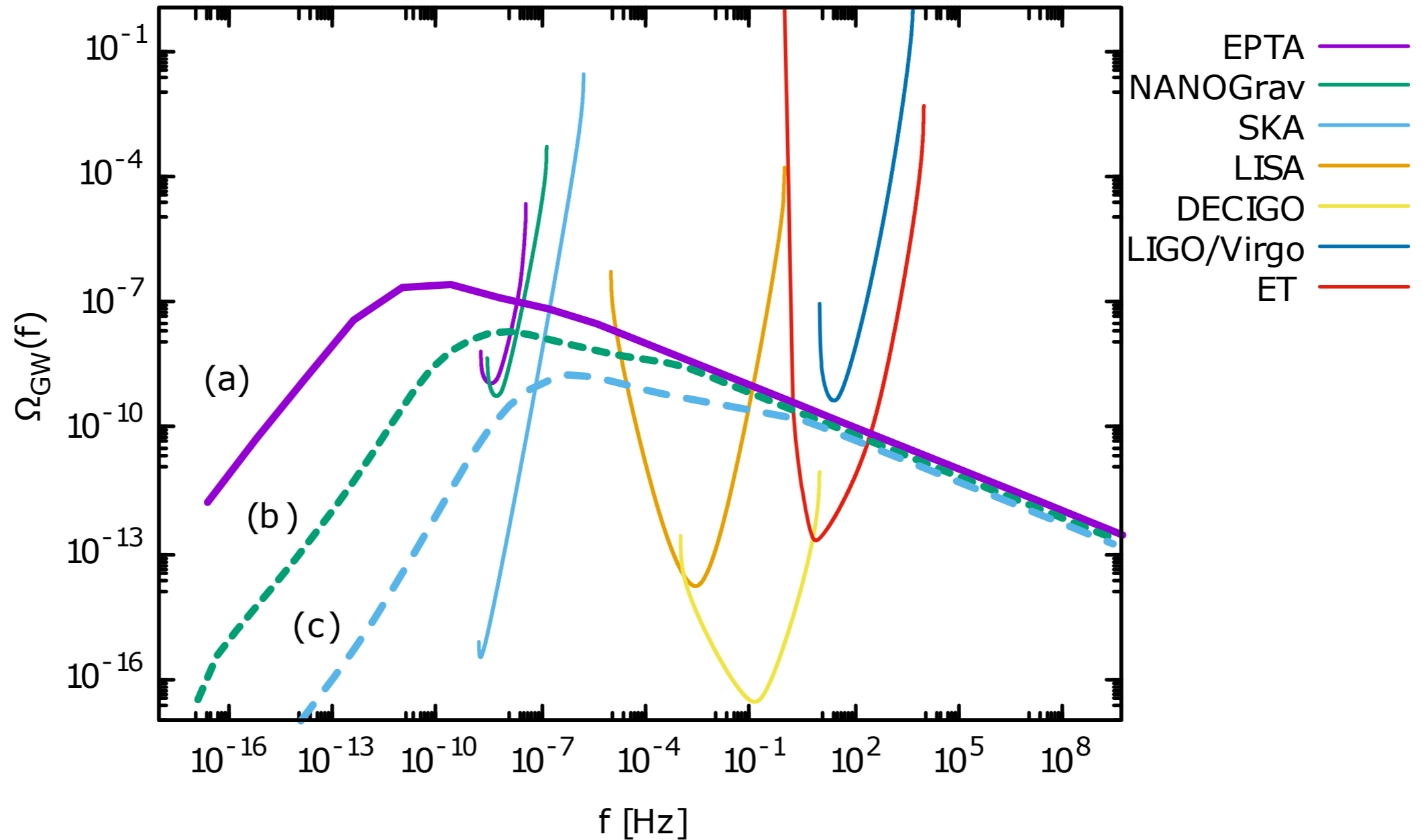


Chang and Cui, 2106.09746

(see also Gorghetto et al, 2101.11007)

GW spectrum

NK, Nakayama 2212.13573



(a) $v = 10^{15}$ GeV, $m_A = 10^{-14}$ eV

(b) $v = 10^{13}$ GeV, $m_A = 10^{-10}$ eV

(c) $v = 10^{12}$ GeV, $m_A = 10^{-5}$ eV

Discussion

More precise study is necessary for

- Scaling violation
- Time-dependence of the tension
- Loop and dark photon production rate
especially near the transition : global \rightarrow local
- Initial loop size distribution (monochromatic or extended?)
- Spectral function of GW from individual loop (cusp- or kink-like?)
(because it is crucial for high frequency region)
- Loop lifetime (deviation from Nambu-Goto string)

discussed in Hindmarsh et al (2017), Matsunami et al (2019)

Summary

- Light dark photon DM can be produced by
 - axion oscillation
 - misalignment mechanism (still viable)
 - decay of cosmic string loops
- Gravitational waves can be a signature of this scenario
 - circular polarization
 - mildly tilted spectrum