

Introduction



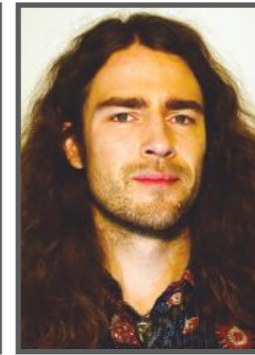
T. Konstandin



H. Rubira



B. Shakya



I. Stomberg



J.v.d.Vis

GWs from
sound waves
in FOPT

Gravitational waves from feebly-interacting particles in a first-order phase transition

Ryusuke Jinno (RESCEU, UTokyo)

2023/11/9, "Gravitational Wave Probes of Physics Beyond Standard Model", Osaka City University

GWs from
FIPs
in FOPT

Summary



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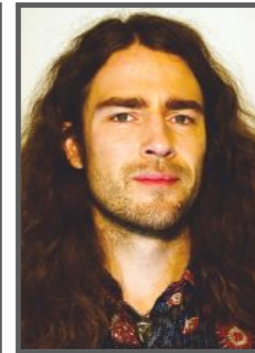
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OVERVIEW OF FIRST-ORDER PHASE TRANSITION & BUBBLE DYNAMICS

[Hogan '83] [Witten '84]

microphysics

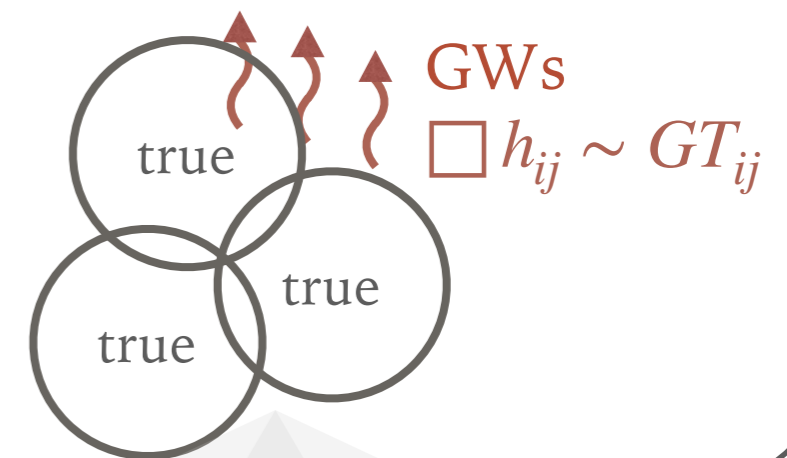
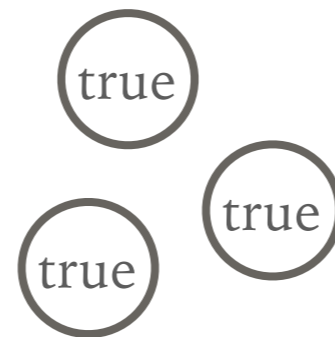
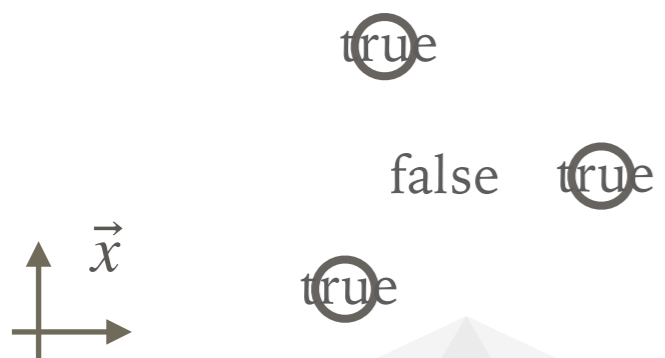
macrophysics

time or scale →

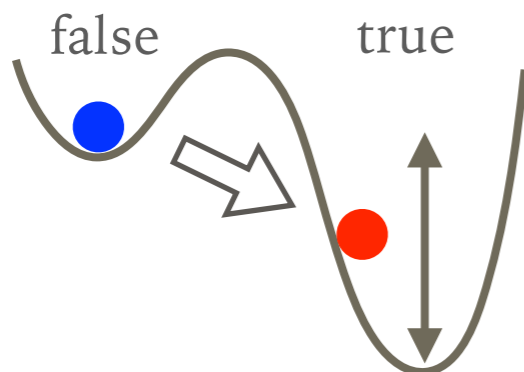
(1) nucleation (核生成)

(2) expansion (拡大)

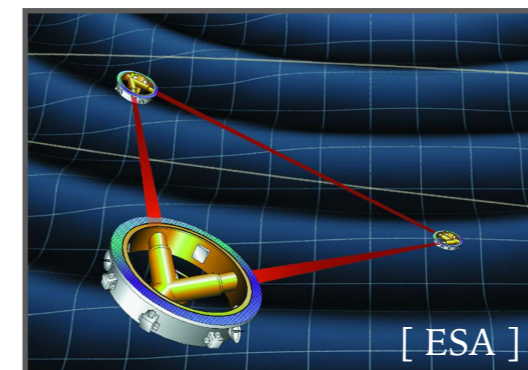
(3) collision (衝突)



Physics of the Higgs sector



GW observations



GW PRODUCTION: THE STANDARD LORE & BEYOND

➤ **GW sources** e.g. [Caprini et al. '16] [Caprini et al. '20]

Bubble walls [Kosowsky, Turner, Watkins '92] [Kosowsky, Turner '92] ...

Energy released accumulates in the walls (= scalar field kinetic & gradient)

Fluid = Sound waves & Turbulence [Kamionkowski, Kosowsky, Turner '93] ...
[Hindmarsh, Huber, Rummukainen, Weir '14] ...

Particles in the broken phase frequently interact and can be described by fluid picture.

Aren't we missing one possibility?

GW PRODUCTION: THE STANDARD LORE & BEYOND

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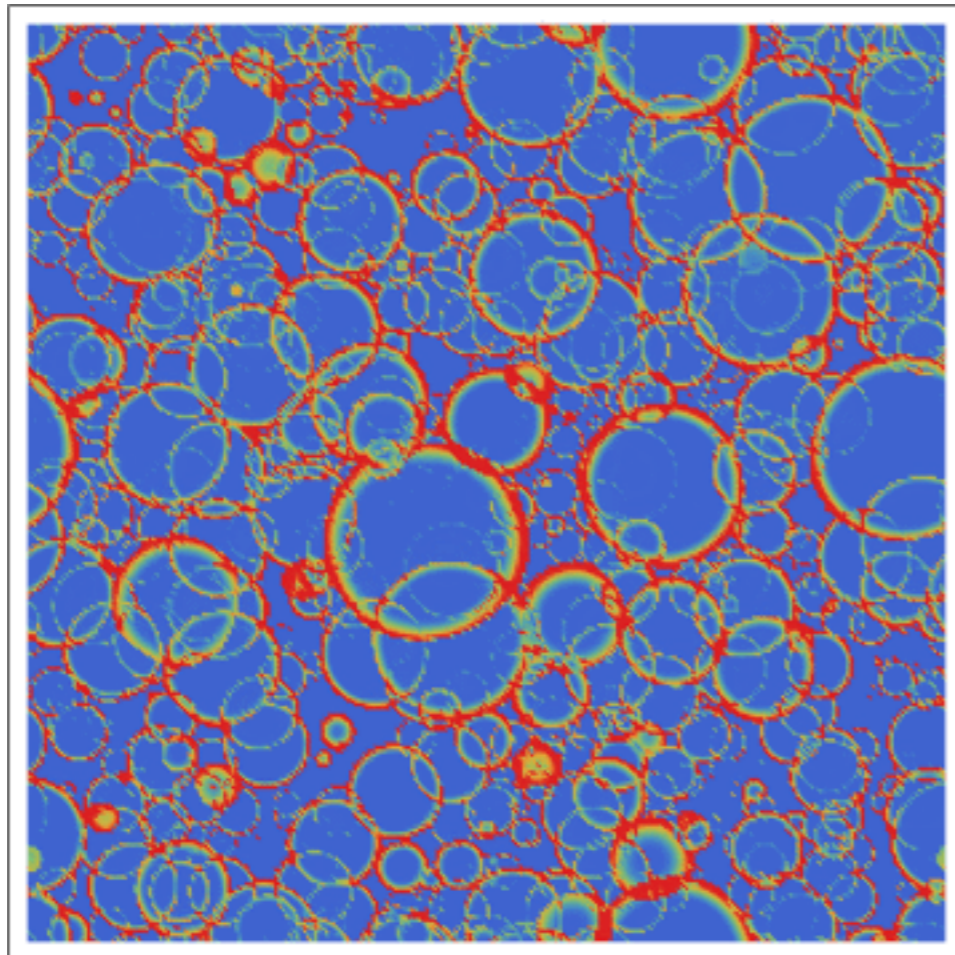
Feebly-interacting particles

Particles in the broken phase are only feebly interacting and free-stream.

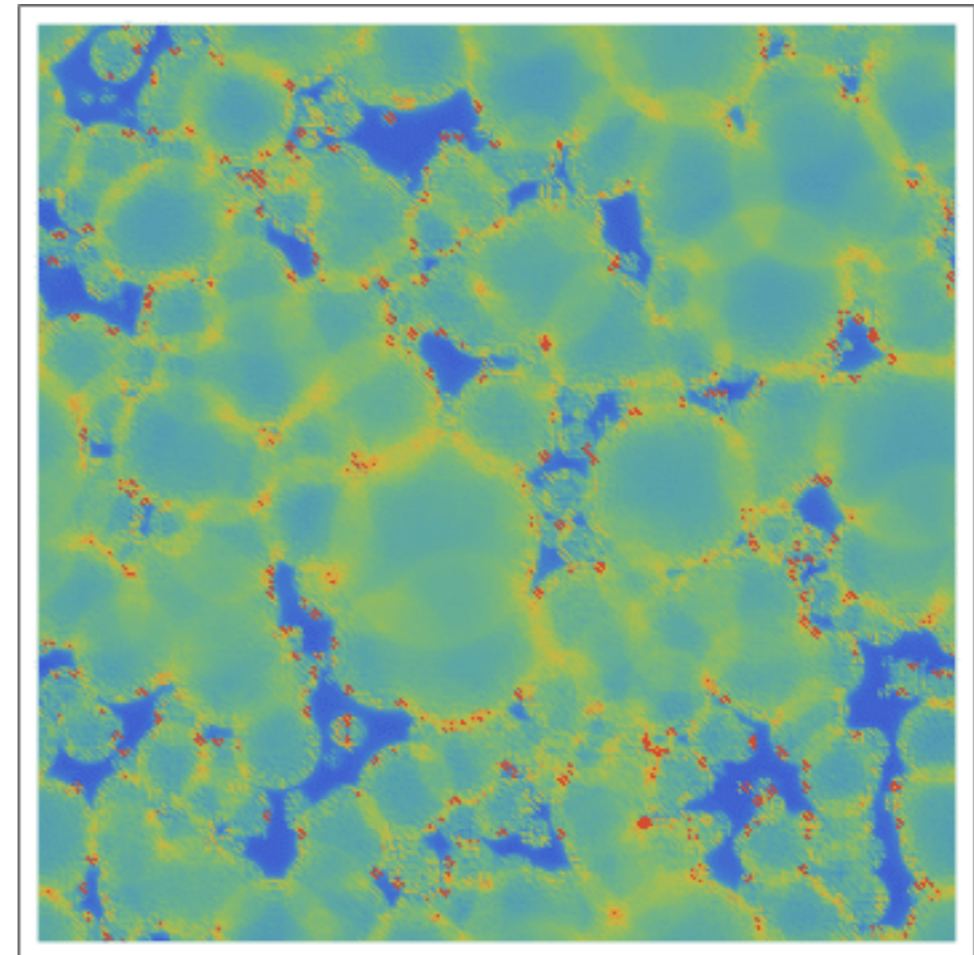
EVOLUTION OF BUBBLES WITH FEEBLY-INTERACTING PARTICLES

- Fluid case vs. feebly-interacting case

Fluid



Feebly-interacting





Introduction

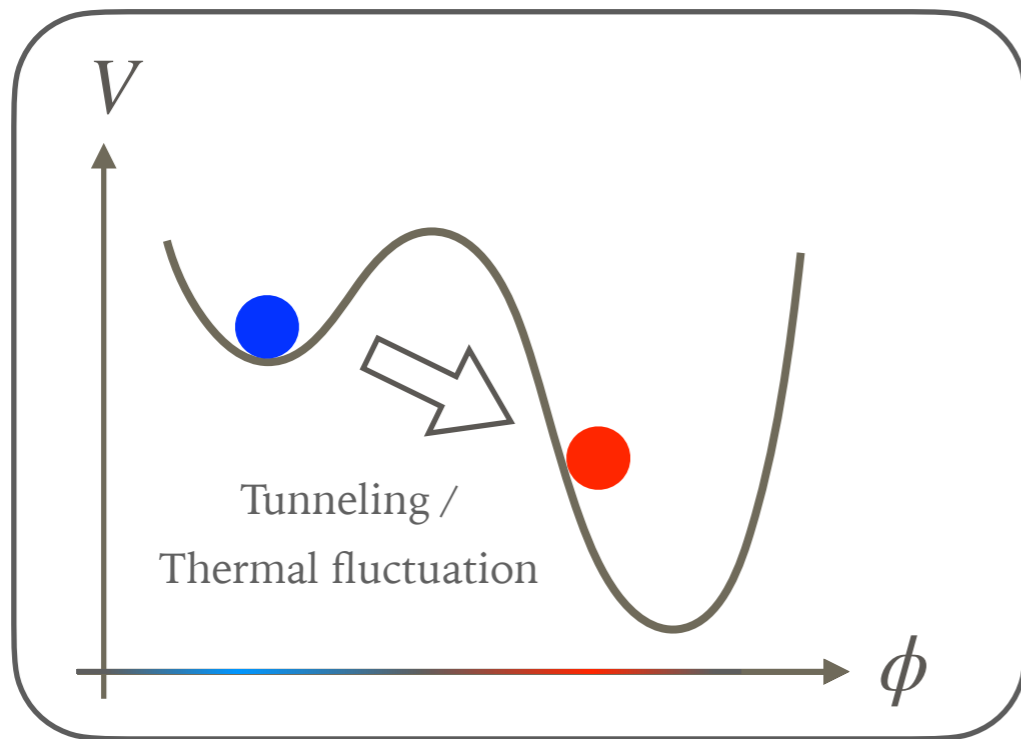
*GWs from
sound waves
in FOPT*

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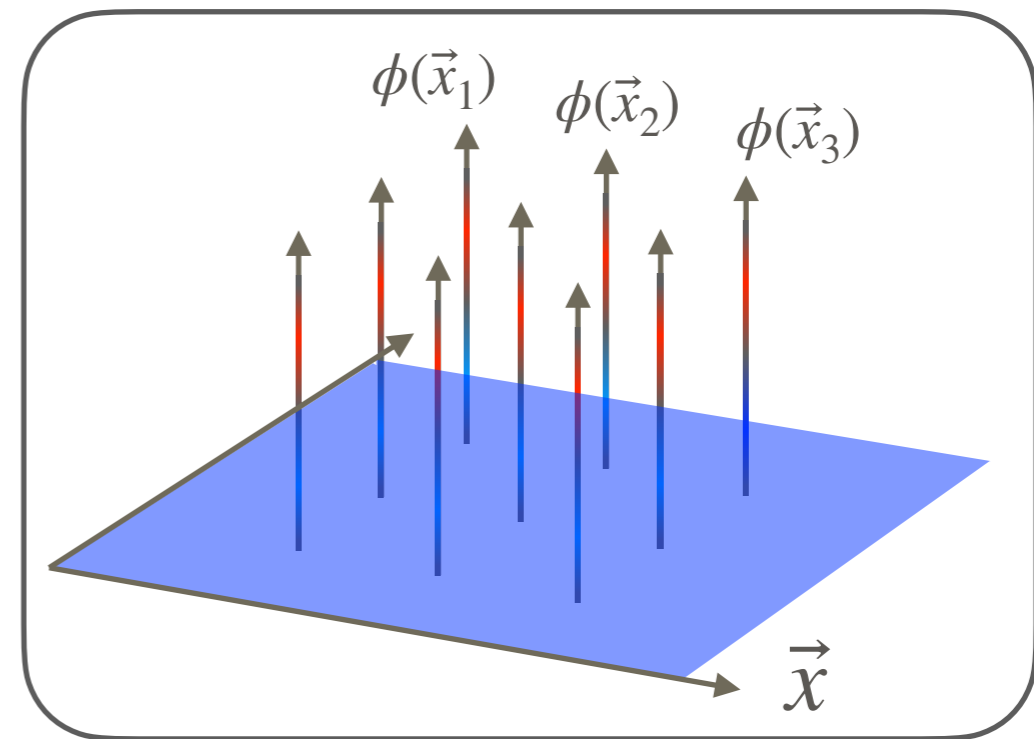
Summary

TUNNELING IN QUANTUM MECHANICS AND QFT

Field space

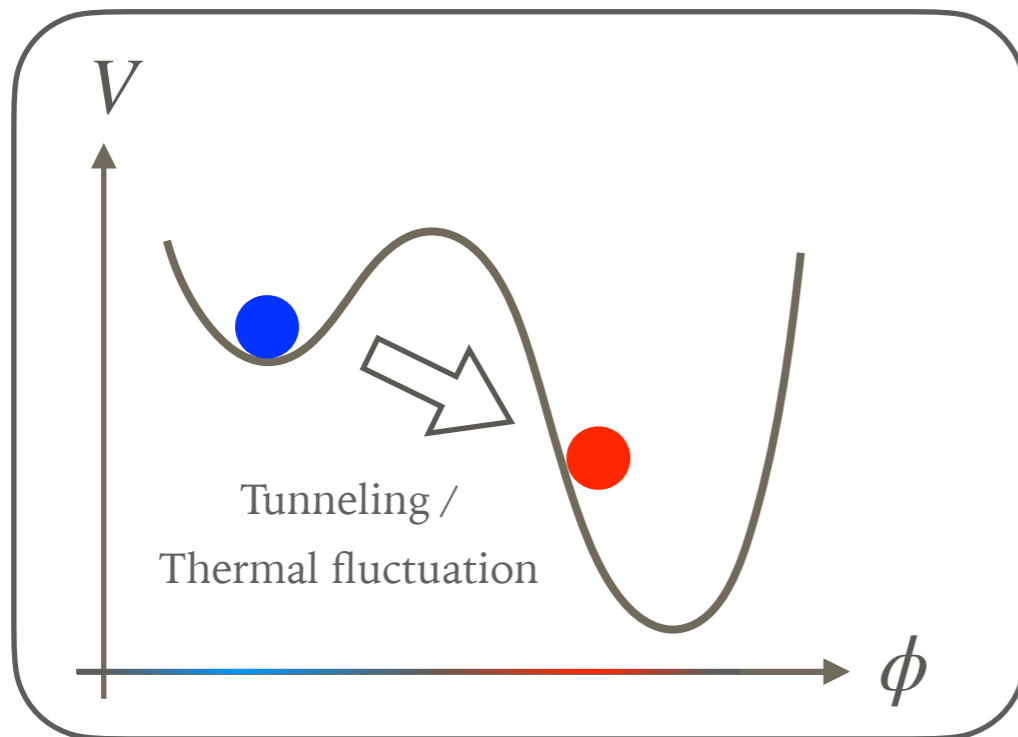


Position space

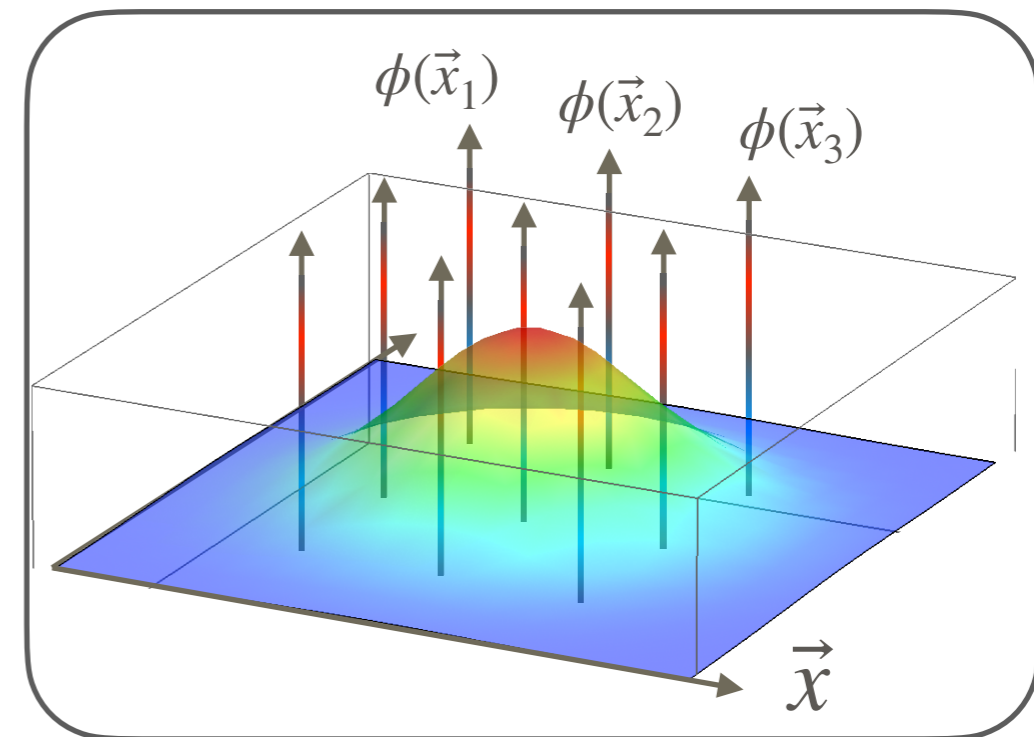


TUNNELING IN QUANTUM MECHANICS AND QFT

Field space

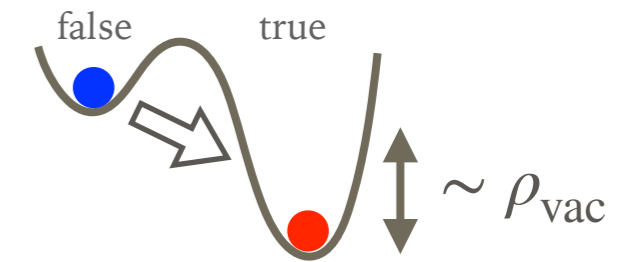


Position space



nucleation (核生成)

STEP 2 : BUBBLE EXPANSION



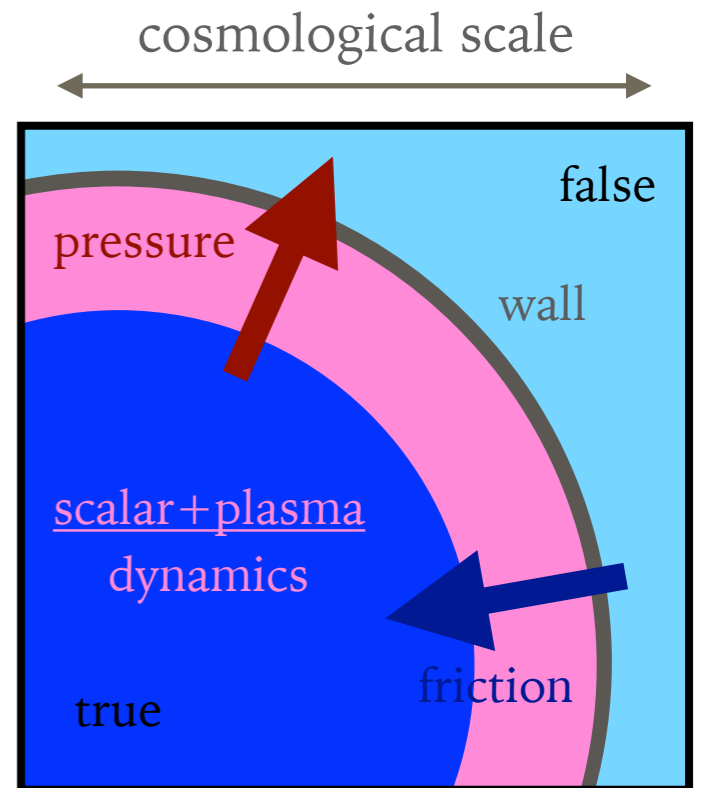
➤ "Pressure vs. Friction" determines the behavior:

(1) **Pressure**: wall is pushed by the released energy

Determined by $\alpha \equiv \rho_{\text{vac}}/\rho_{\text{plasma}}$

see e.g. [Espinosa et al. '10, Hindmarsh et al. '15, Giese et al. '20, Giese et al. '21]

(2) **Friction**: wall is pushed back by plasma particles



STEP 2 : BUBBLE EXPANSION

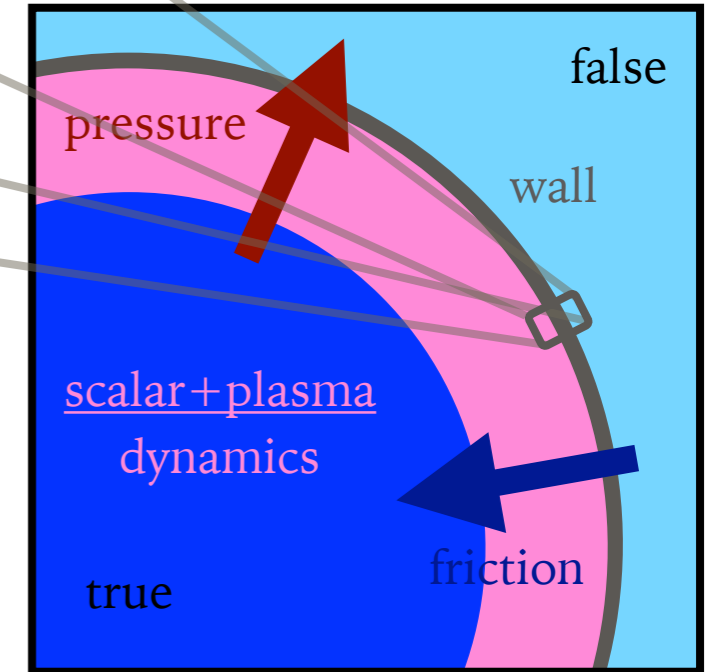
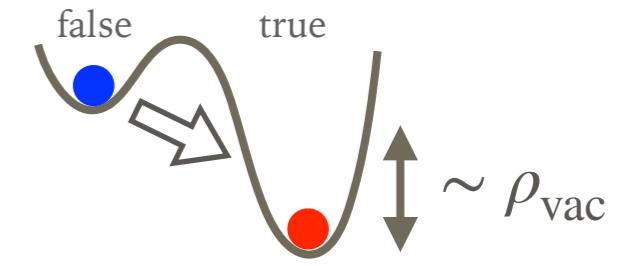
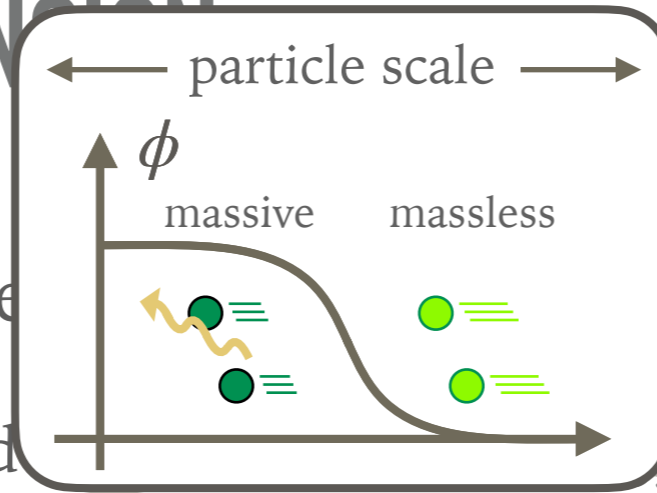
► "Pressure vs. Friction" determines

(1) Pressure: wall is pushed

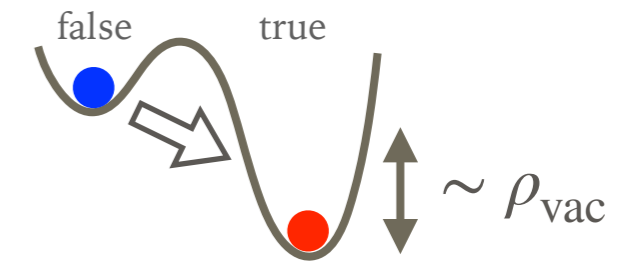
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(2) Friction: wall is pushed back by plasma particles



STEP 2 : BUBBLE EXPANSION



➤ "Pressure vs. Friction" determines the behavior:

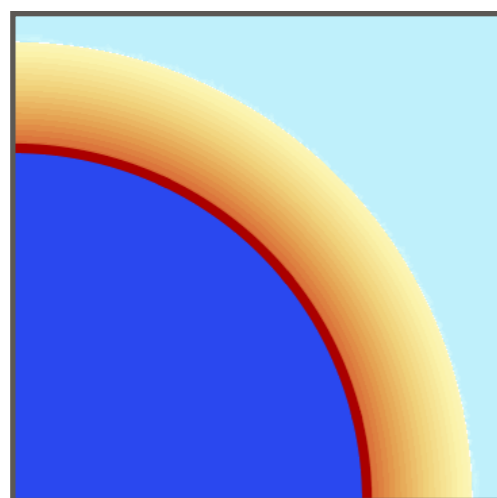
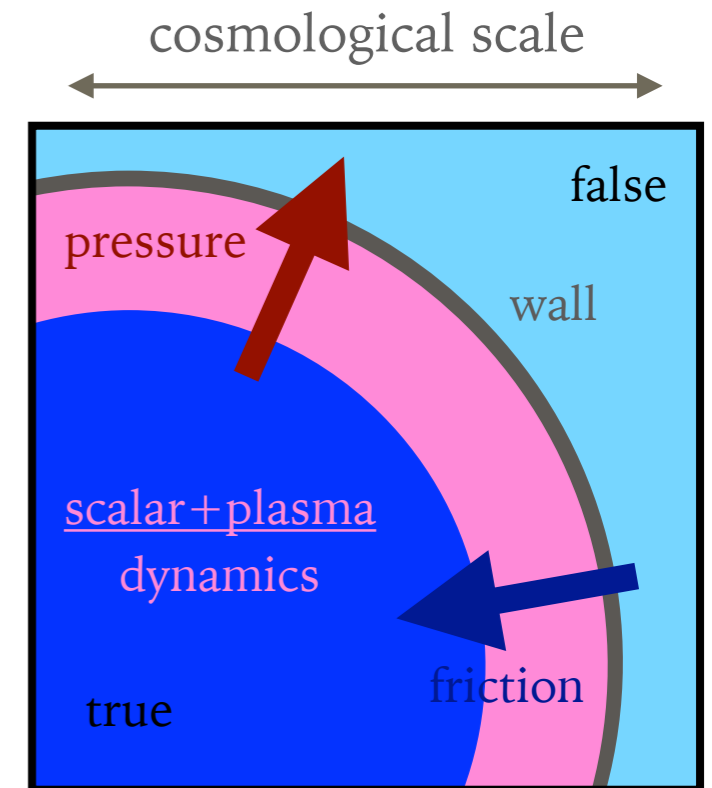
(1) **Pressure**: wall is pushed by the released energy

Determined by $\alpha \equiv \rho_{\text{vac}}/\rho_{\text{plasma}}$

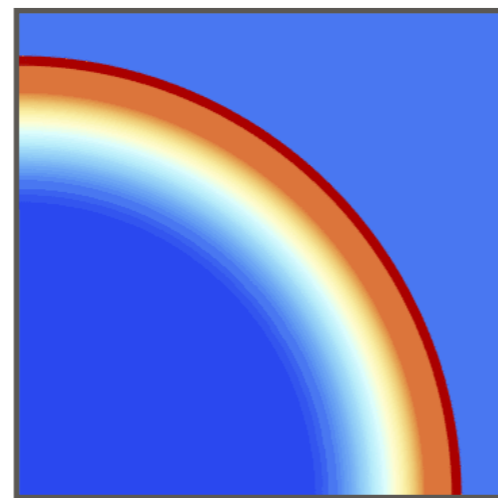
see e.g. [Espinosa et al. '10, Hindmarsh et al. '15, Giese et al. '20, Giese et al. '21]

(2) **Friction**: wall is pushed back by plasma particles

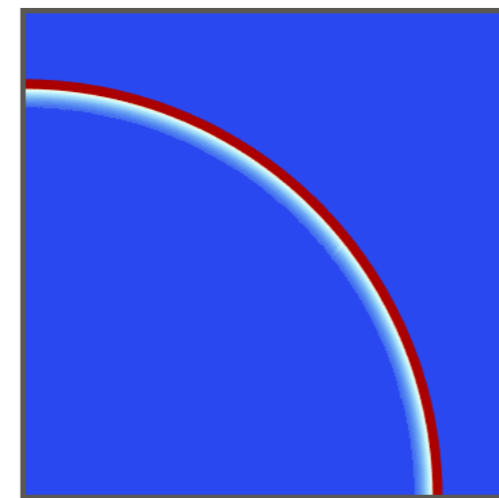
➤ Different types of bubble expansion



deflagration



detonation



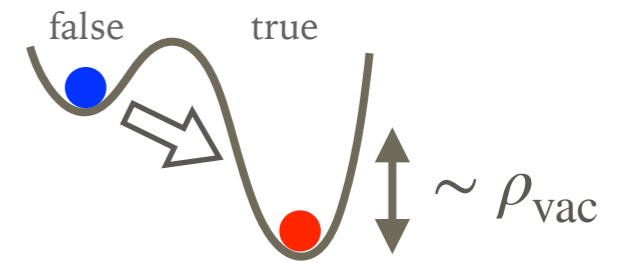
~ 1 relativistic detonation $\gg 1$



runaway



STEP 2 : BUBBLE EXPANSION

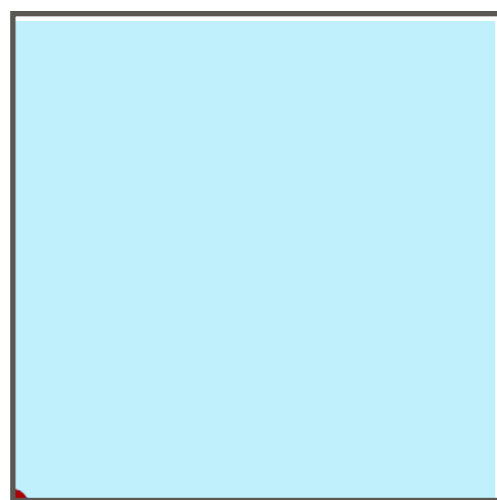
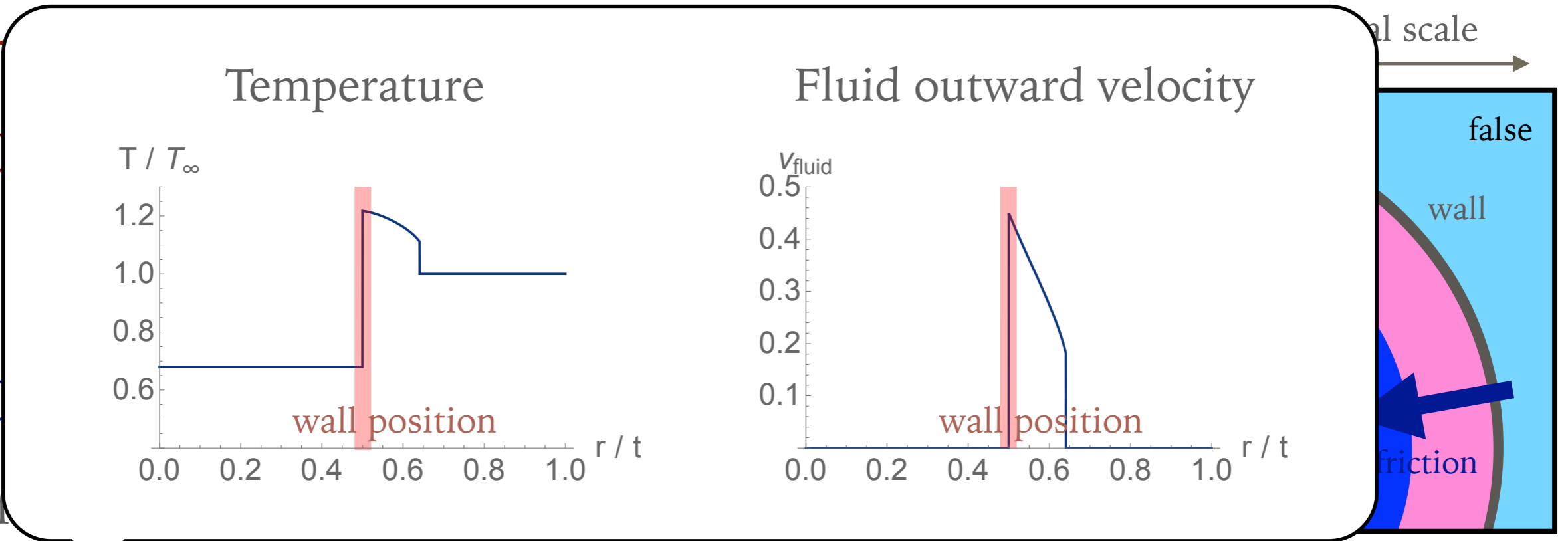


➤ "Pr"

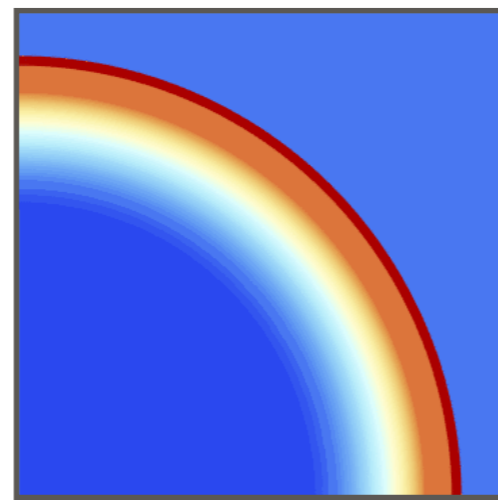
(1)

(2)

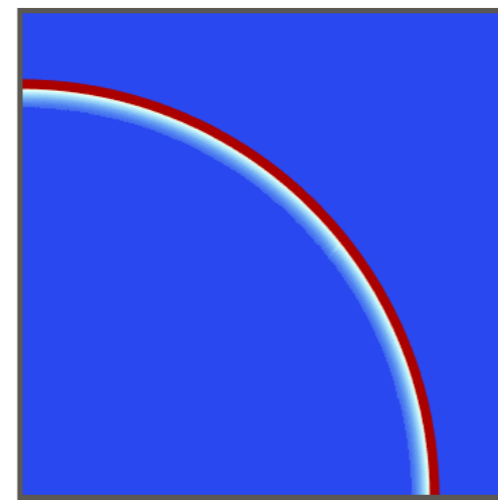
➤ Dis



deflagration



detonation



~ 1 relativistic detonation $\gg 1$

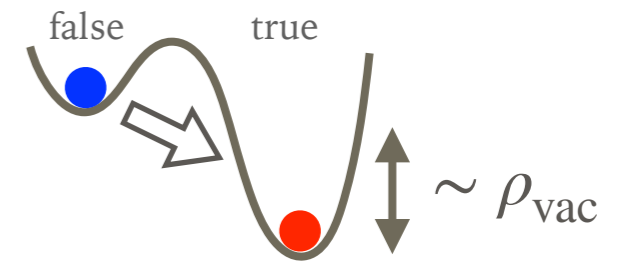


runaway

————— α —————



STEP 2 : BUBBLE EXPANSION

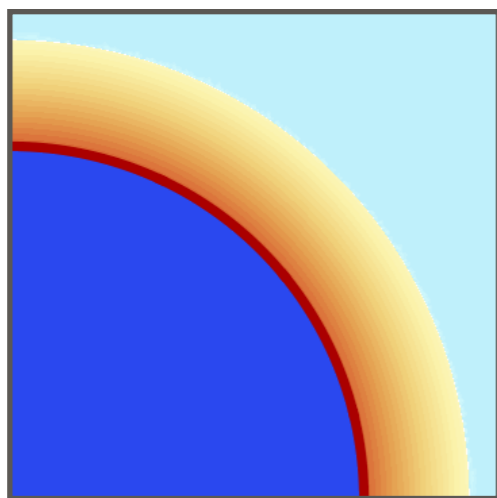
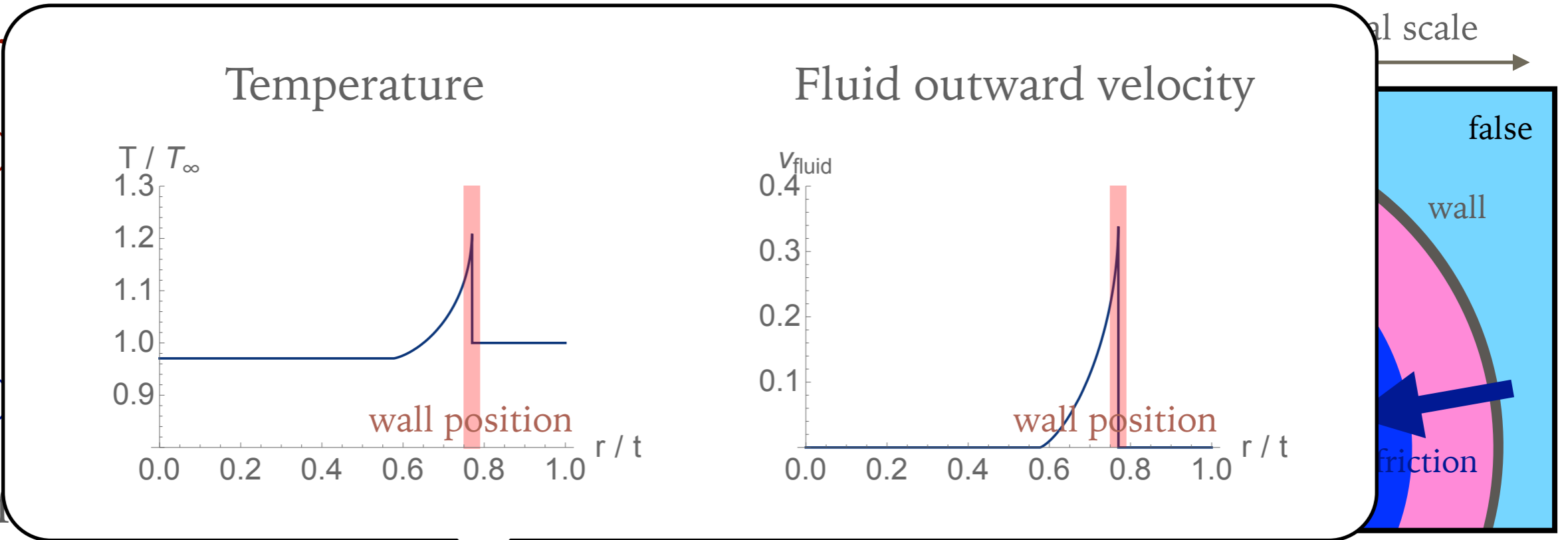


➤ "Pr"

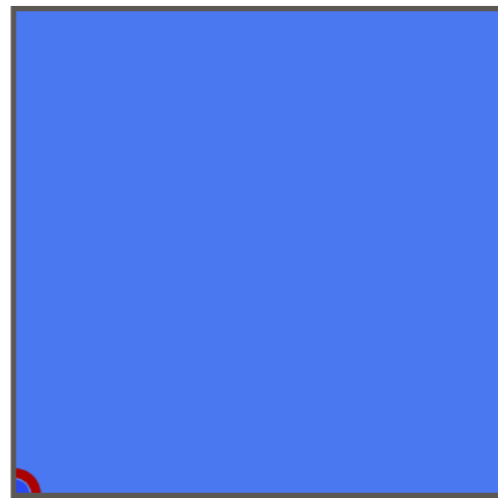
(1)

(2)

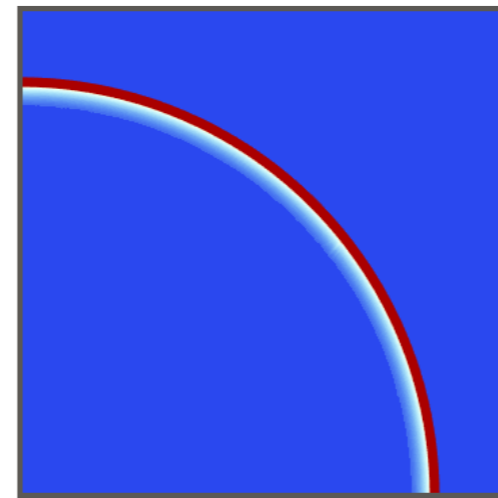
➤ Dis



deflagration



detonation



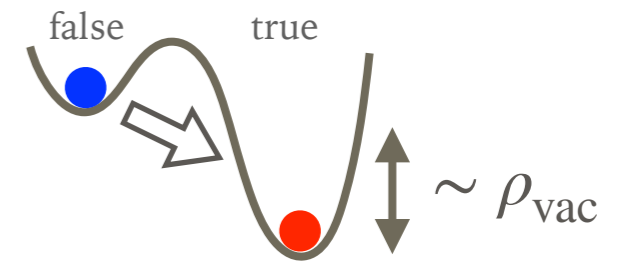
~ 1 relativistic detonation $\gg 1$



runaway



STEP 2 : BUBBLE EXPANSION

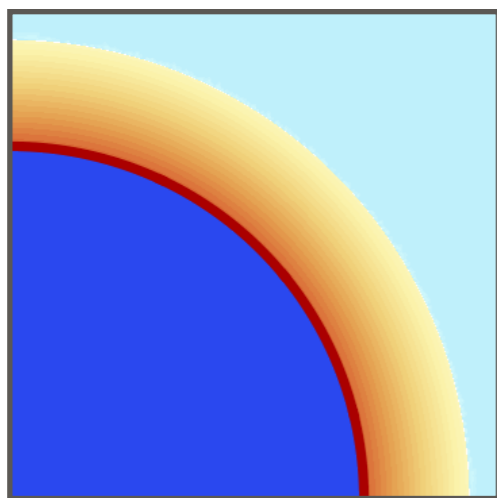
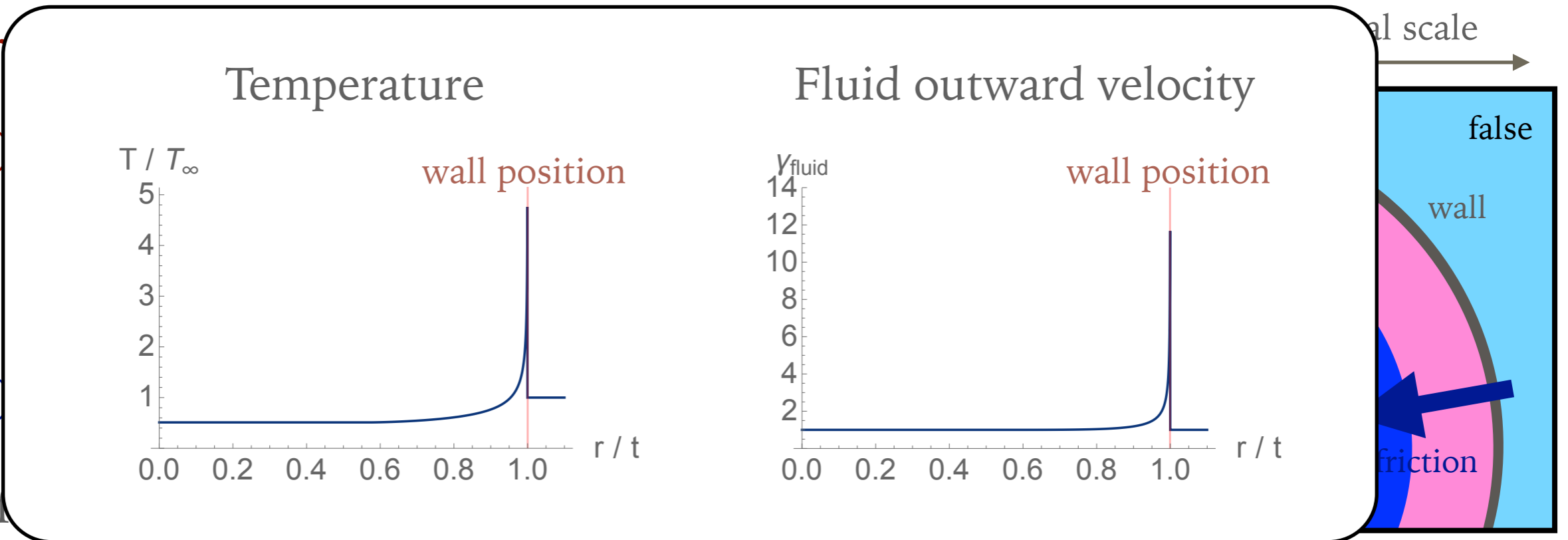


➤ "Pr

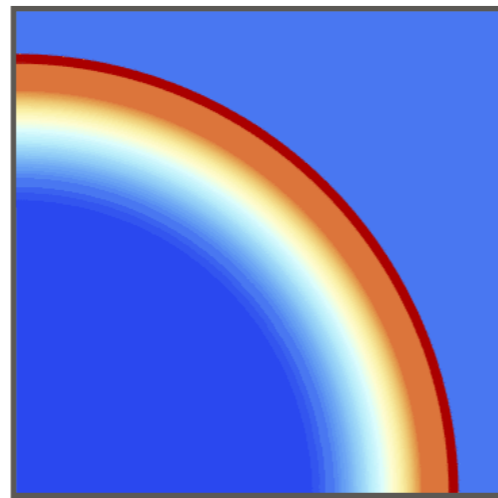
(1)

(2)

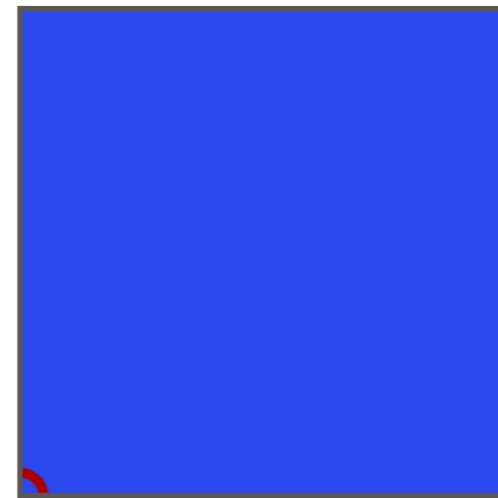
➤ Dis



deflagration



detonation



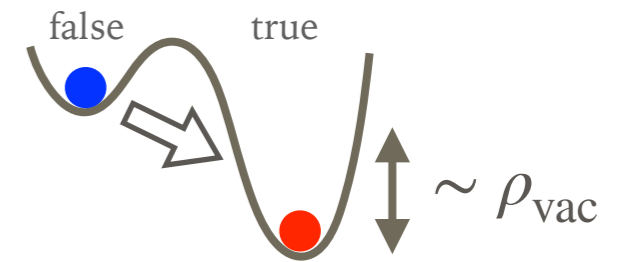
~ 1 relativistic detonation $\gg 1$



runaway



STEP 2 : BUBBLE EXPANSION



➤ "Pr

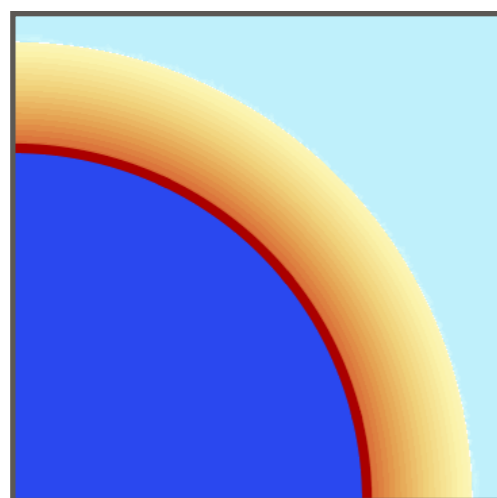
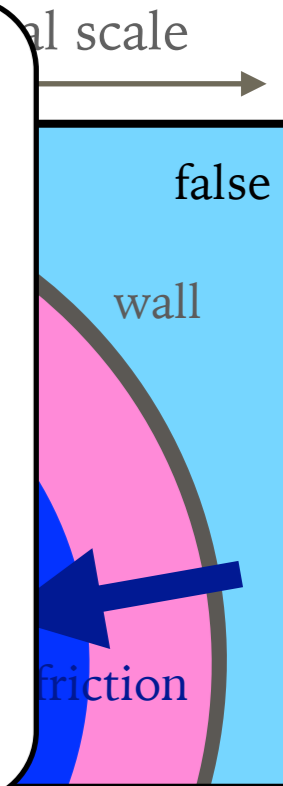
(1)

Plasma particles cannot stop the acceleration of the walls:
walls continue to accelerate until they collide with others

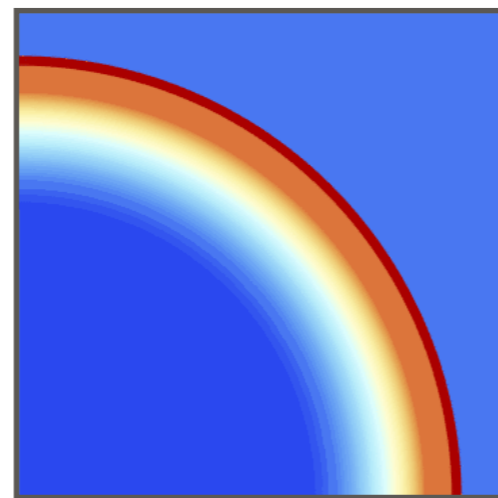
(2)

[Bodeker & Moore '09, '17] [Azatov & Vanvlasselaer '20] [Gouttenoire, RJ, Sala '21]
[Höche, Kozaczuk, Long, Turner, Wang '20] [Barroso Mancha, Prokopec, Świeżewska '21]
[Azatov, Barni, Petrossian-Byrne, Vanvlasselaer '23] ...

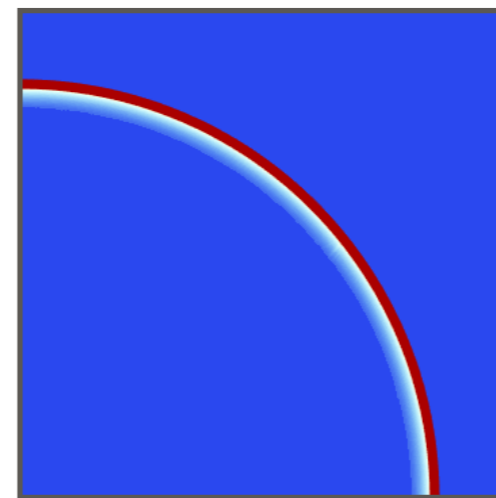
➤ Dis



deflagration



detonation



~ 1 relativistic detonation $\gg 1$



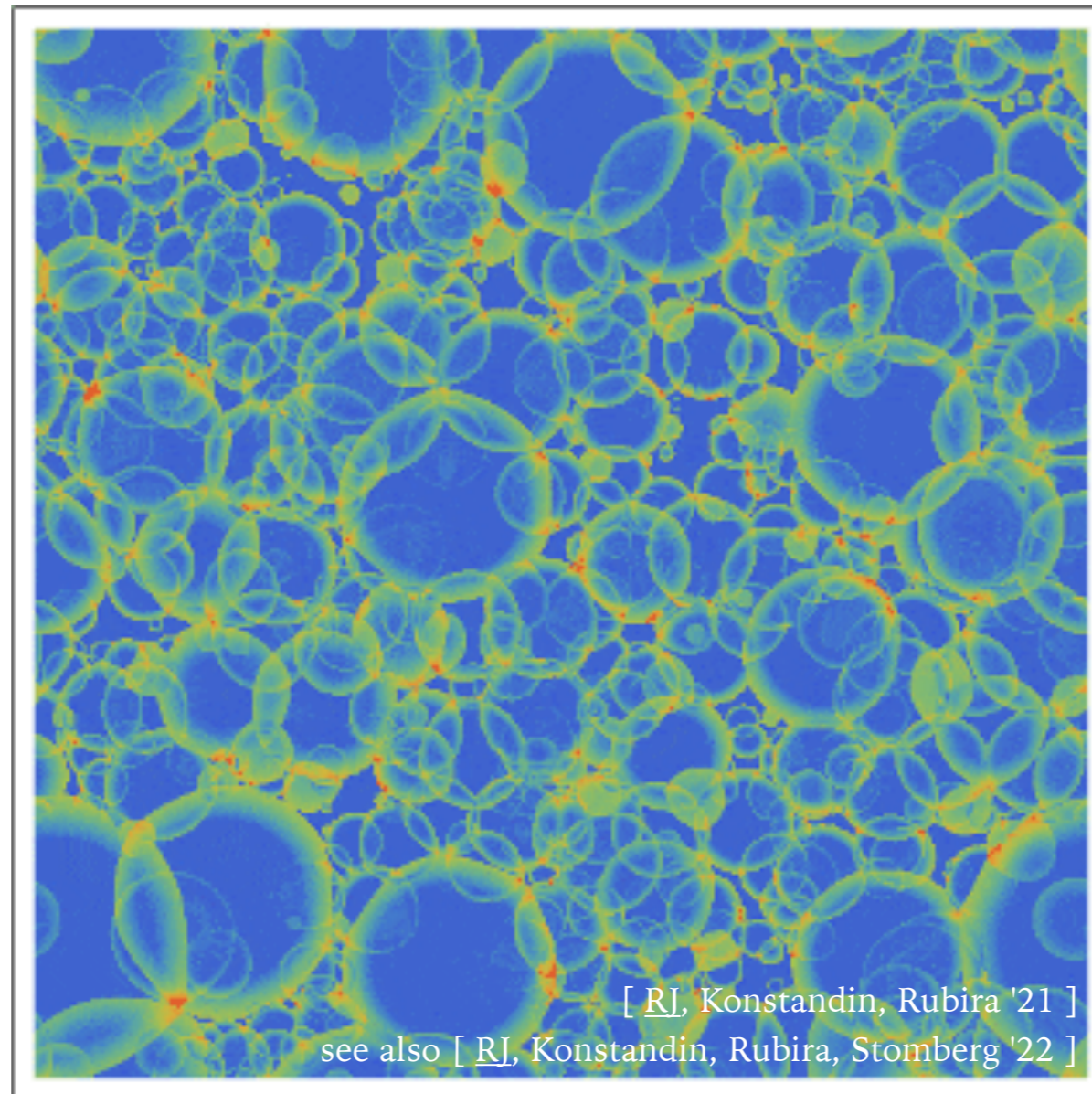
runaway



STEP 3: BUBBLE COLLISION & FLUID DYNAMICS

- Bubbles collide, and fluid dynamics sets in (example for )

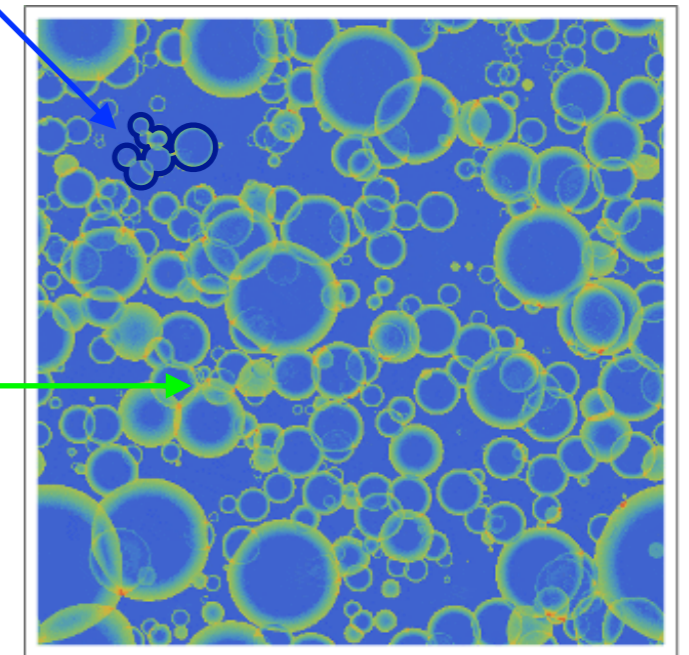
[Hindmarsh, Huber, Rummukainen, Weir '14, '15, '17] [Hindmarsh '15, +Hijazi '19]



GRAVITATIONAL WAVE SOURCES

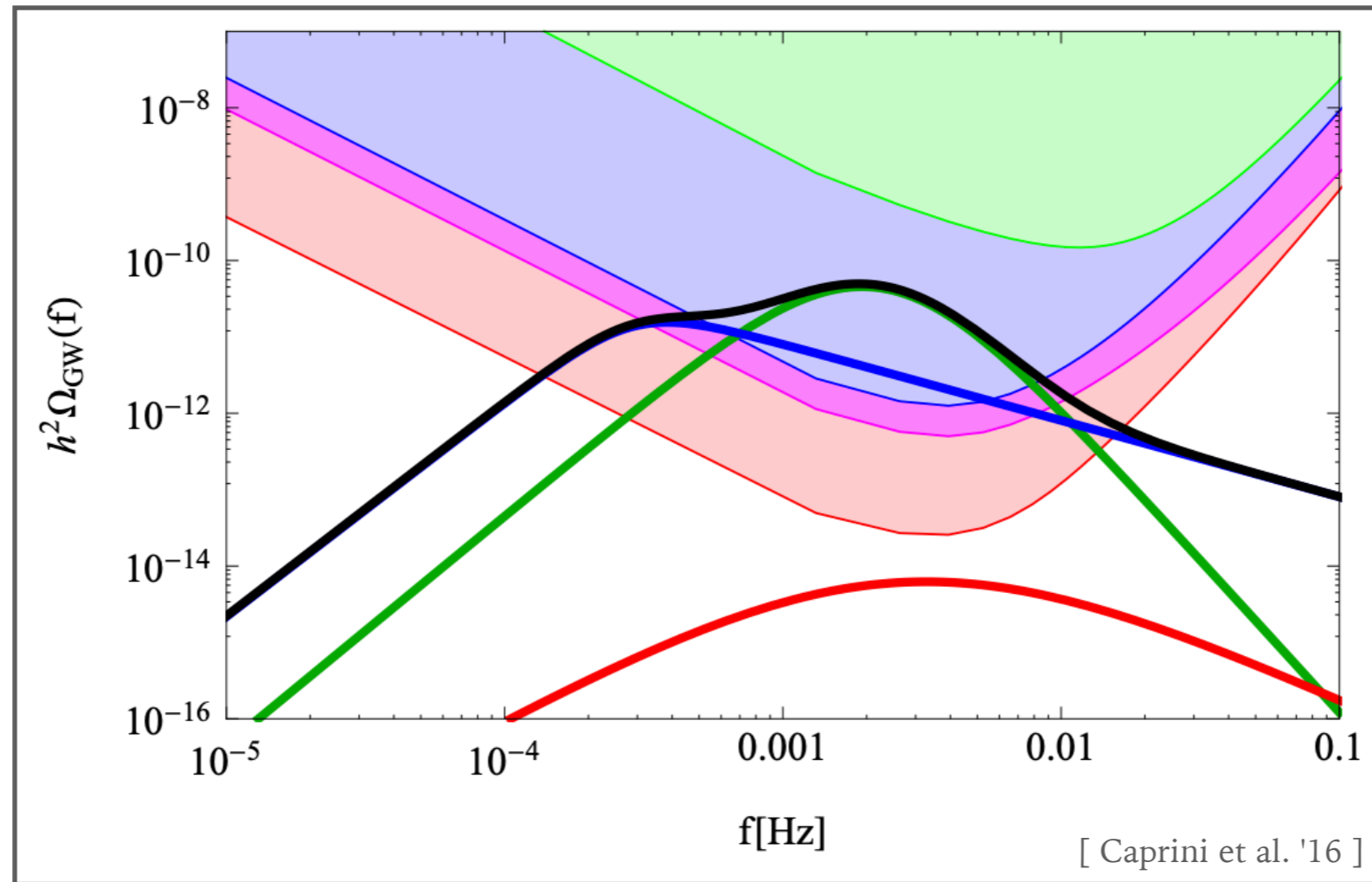
see e.g. [Caprini et al. '16] [Caprini et al. '20]

- ▶ Bubble collision
 - Kinetic & gradient energy of the scalar field (= order parameter field)
 - Dominant when the transition is extremely strong and the walls runaway
- ▶ Sound waves
 - Compression mode of the fluid motion
 - Dominant unless the transition is extremely strong
- ▶ Turbulence
 - Turbulent motion caused by fluid nonlinearity
 - Expected to develop at a later stage



important at later stage

GRAVITATIONAL WAVE SPECTRUM





Introduction

*GWs from
sound waves
in FOPT*

*GWs from
FIPs
in FOPT*

Summary

GW PRODUCTION: THE STANDARD LORE & BEYOND

➤ GW sources e.g. [Caprini et al. '16] [Caprini et al. '20]

Bubble walls [Kosowsky, Turner, Watkins '92] [Kosowsky, Turner '92] ...

Energy released accumulates in the walls (= scalar field kinetic & gradient)

Fluid = Sound waves & Turbulence [Kamionkowski, Kosowsky, Turner '93] ...
[Hindmarsh, Huber, Rummukainen, Weir '14] ...

Particles in the broken phase frequently interact and can be described by fluid picture.

Feebly-interacting particles

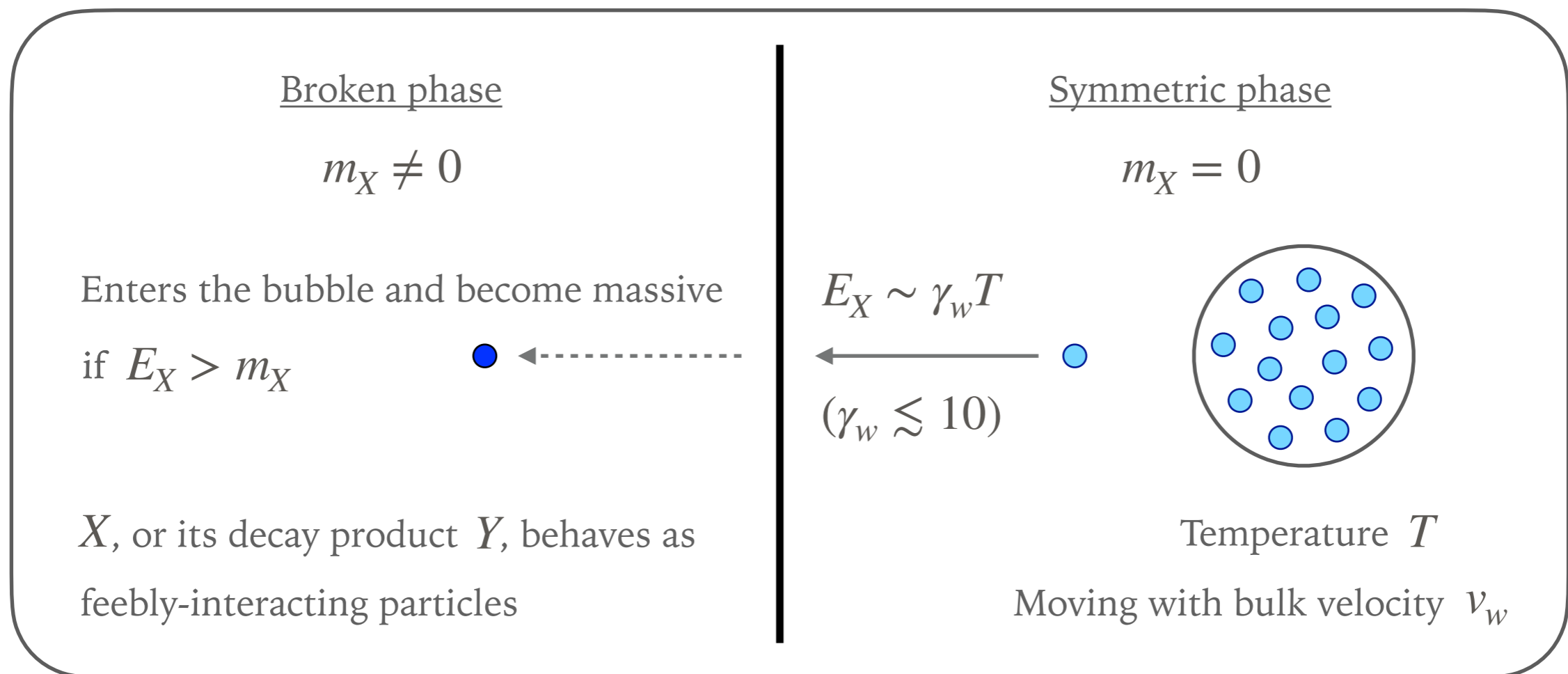
Particles in the broken phase are only feebly interacting and free-stream.

PARTICLE PHYSICS FRAMEWORK

► Setup:

In the broken phase, particles or their decay product free-stream

Wall rest frame



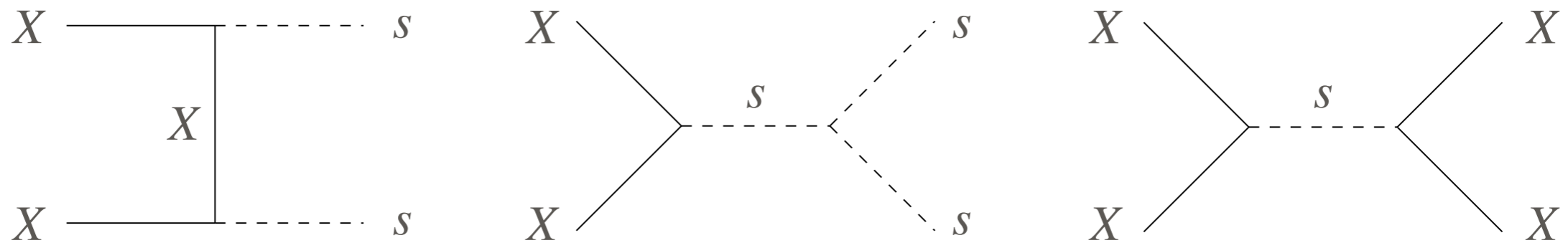
PARTICLE PHYSICS FRAMEWORK

- Consider a dark-sector thermal bath, with temperature T
- ...that undergoes a first-order phase transition
 - scalar field s drives the transition
 - bubble walls reach a terminal velocity v_w (or equivalently $\gamma_w = 1/\sqrt{1 - v_w^2}$)
due to the coupling to particle X
- ...and also produces feebly-interacting particles
 - particle X becomes massive when crossing the wall
 - either X or its decay product behaves as feebly-interacting particles

CONDITIONS ON FEEBLE INTERACTION

► How do X particles interact?

- Couplings that gives rise to mass also give rise to interactions

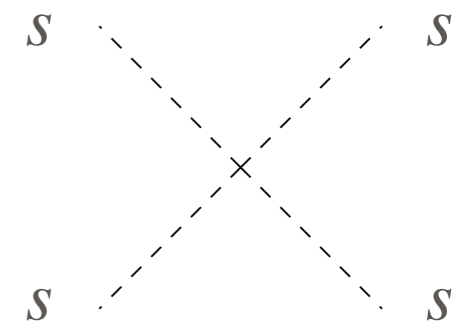


► Can $X = s$ free-stream? → No

- s has to gain a large mass

(for s particles to dominantly contribute to the friction),

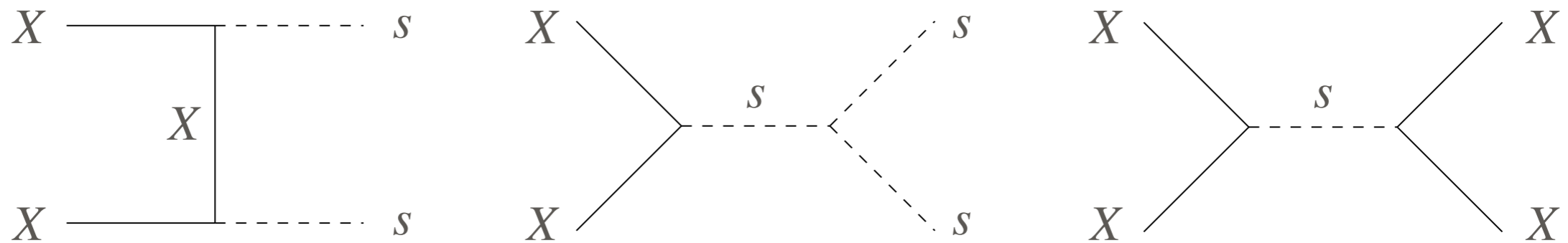
but this immediately means large quartic among s particles



CONDITIONS ON FEEBLE INTERACTION

► How do X particles interact?

- Couplings that gives rise to mass also give rise to interactions



► Can $X = s$ free-stream? → No

► Can $X = Z'$ boson free-stream? → Doable, but not generic

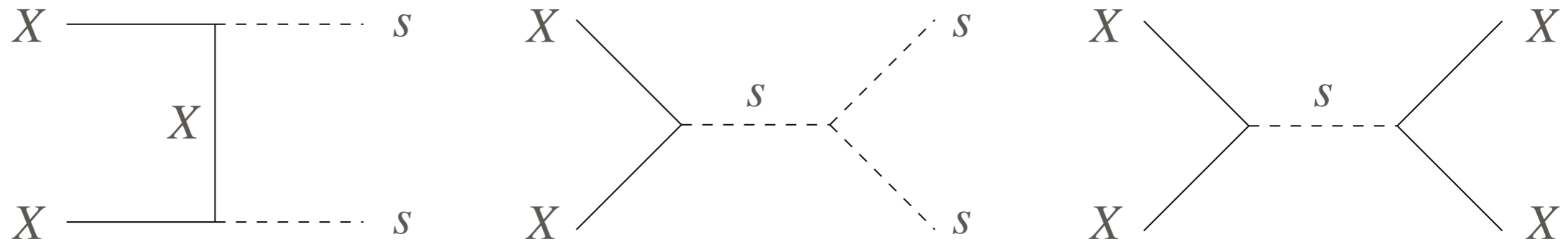
- Condition to free-stream over a typical bubble size R_*

$$n_X \sigma R_* \sim T^3 \frac{g'^4}{(4\pi)^2} \frac{m_{Z'}^2}{m_s^4} \frac{1}{\beta} \sim \frac{g'^6}{\lambda_s^2} \frac{T \overset{\text{large}}{M_P} H}{\langle s \rangle^2 \beta} \lesssim 1$$

CONDITIONS ON FEEBLE INTERACTION

► How do X particles interact?

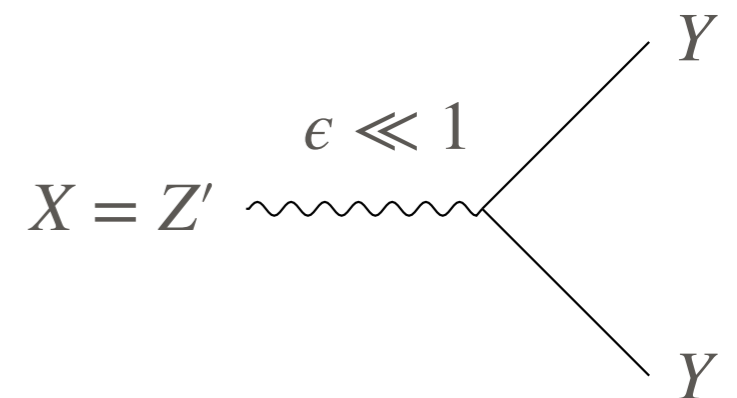
- Couplings that gives rise to mass also give rise to interactions



► Can $X = s$ free-stream? → No

► Can $X = Z'$ boson free-stream? → Doable, but not generic

► Can a decay product of X free-stream? → Yes



SINGLE-BUBBLE PROFILE

- Energy-momentum tensor of a single bubble before collision

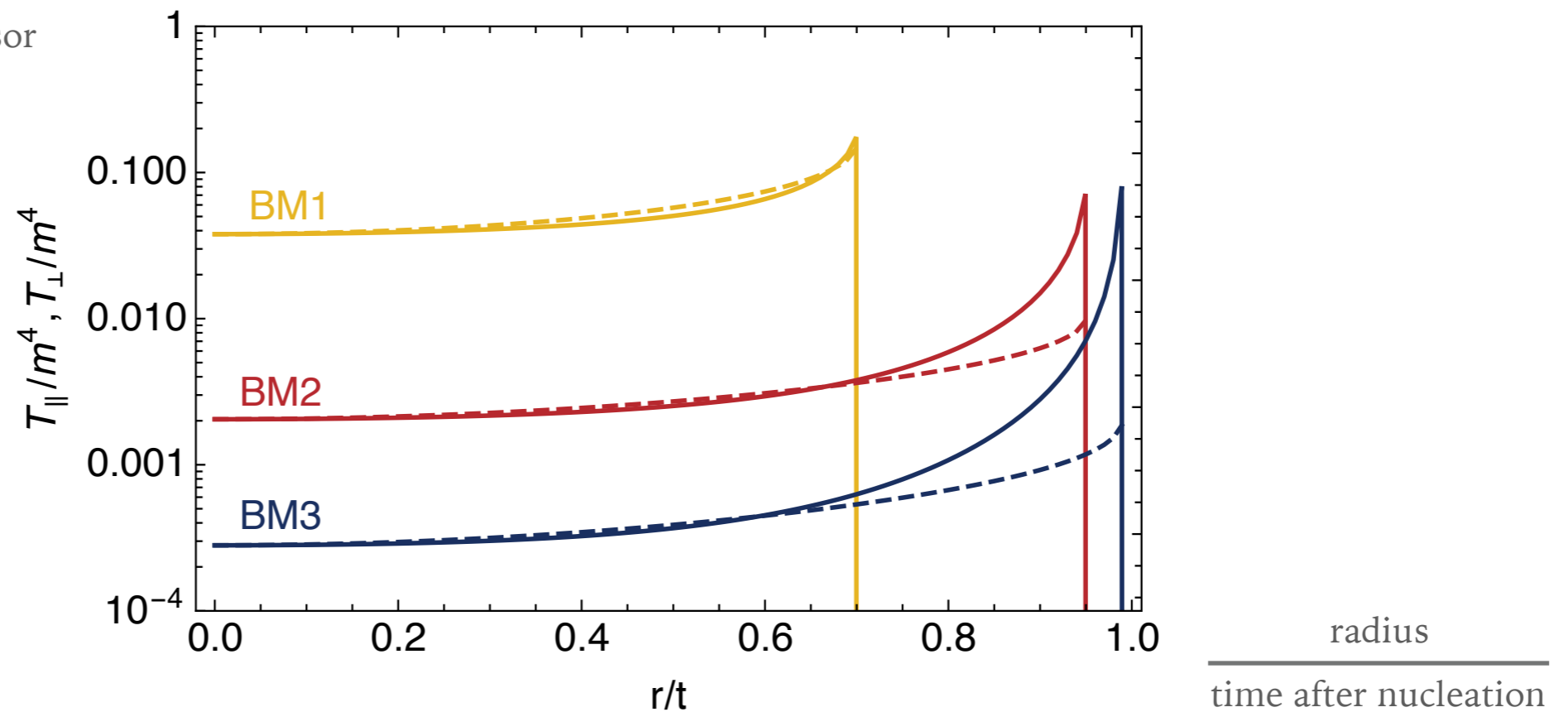
We take 3 benchmark points: $\left\{ \begin{array}{l} m/T = 1, v_w = 0.7 \quad (\text{BM1}) \\ m/T = 2, v_w = 0.95 \quad (\text{BM2}) \\ m/T = 3, v_w = 0.99 \quad (\text{BM3}) \end{array} \right.$

see also [Lewicki, Vaskonen, Veermäe '22]

Energy-momentum tensor

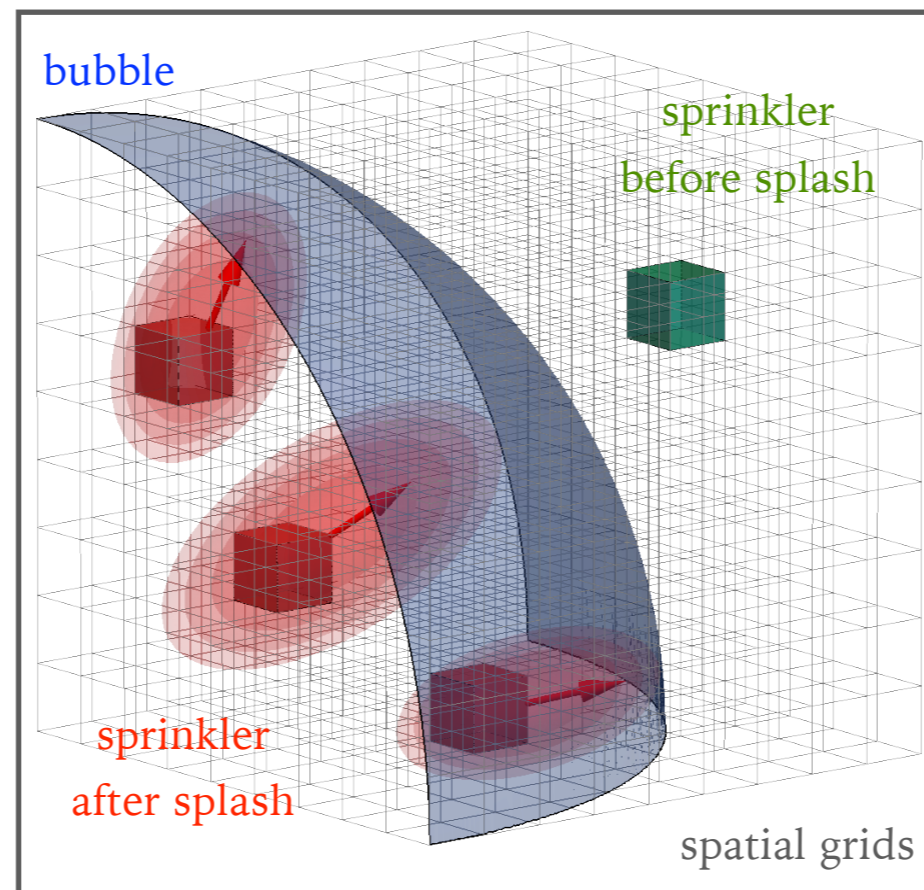
$$T_{\parallel} = T_{zz}$$

$$T_{\perp} = T_{xx} = T_{yy}$$



HOW TO CALCULATE THE GW SPECTRUM

- To calculate the GW spectrum, we don't use the time evolution shown in the animation in the previous slides
- We instead propose a new calculation scheme – "sprinkler picture"



SPRINKLER PICTURE FOR GW CALCULATION

➤ How the "sprinkler picture" works

① Imagine **each grid point has a sprinkler** that splashes free-streaming particles when hit by the wall

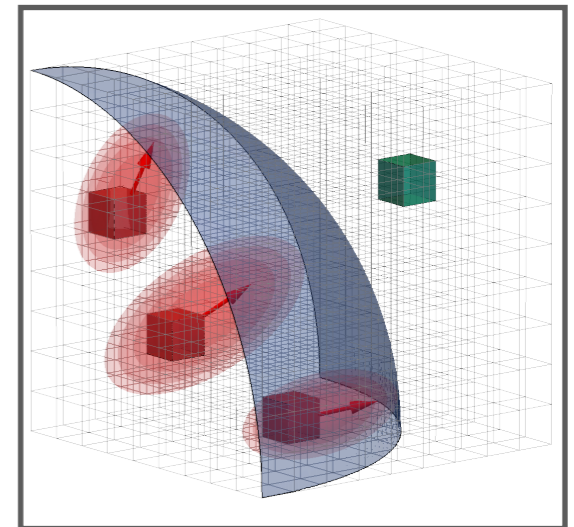
② **The sprinklers are universal:**

their only difference is when and in which direction they are hit

③ GW production from one sprinkler is easily calculable.

Contributions from different sprinklers (= grids) are linearly superposed.

➤ This method is possible because GW production is linear in each sprinkler in the present system

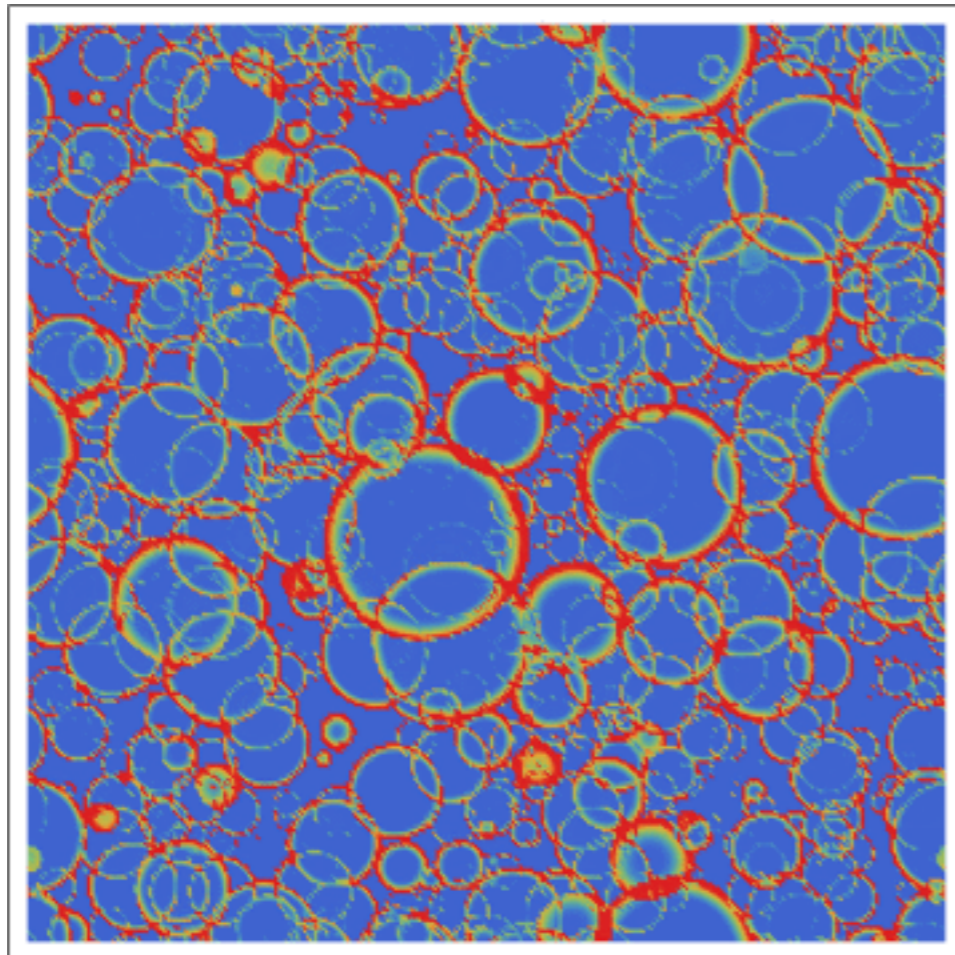


$$\square h_{ij} \sim T_{ij} \sim \sum_{\text{sprinkler } p} T_{ij}^{(p)}$$

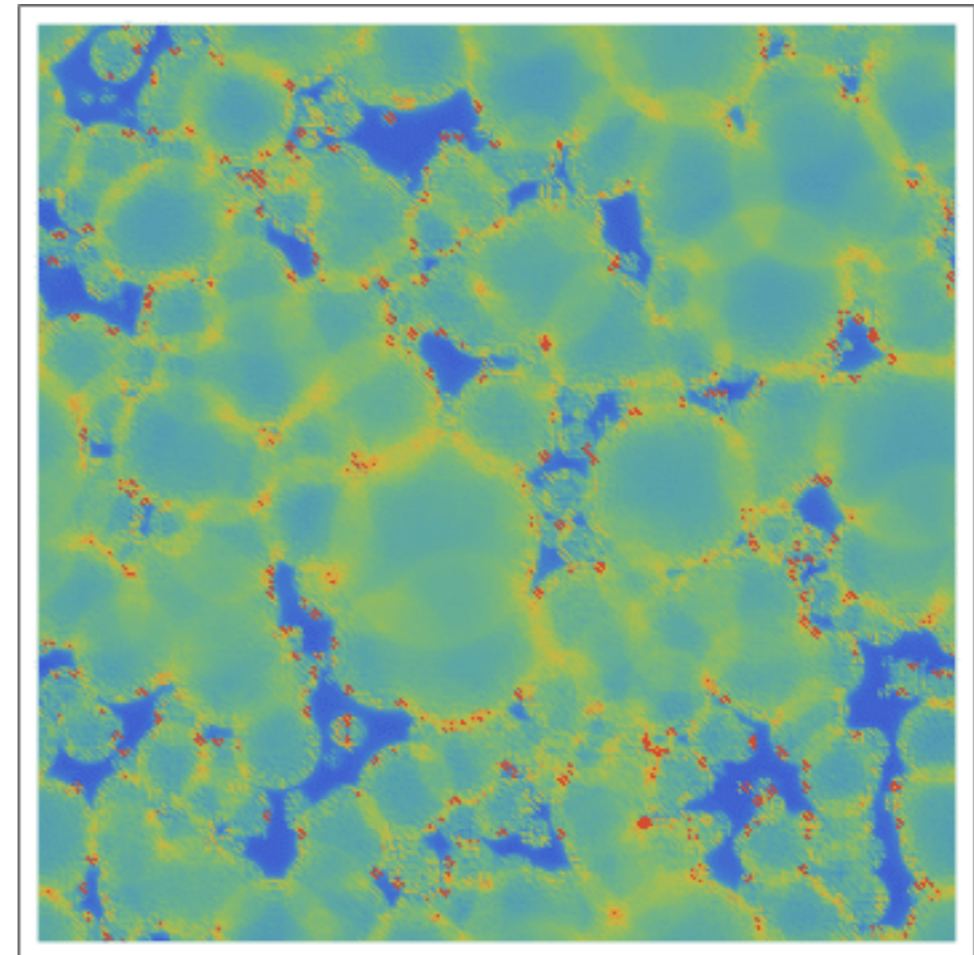
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Fluid



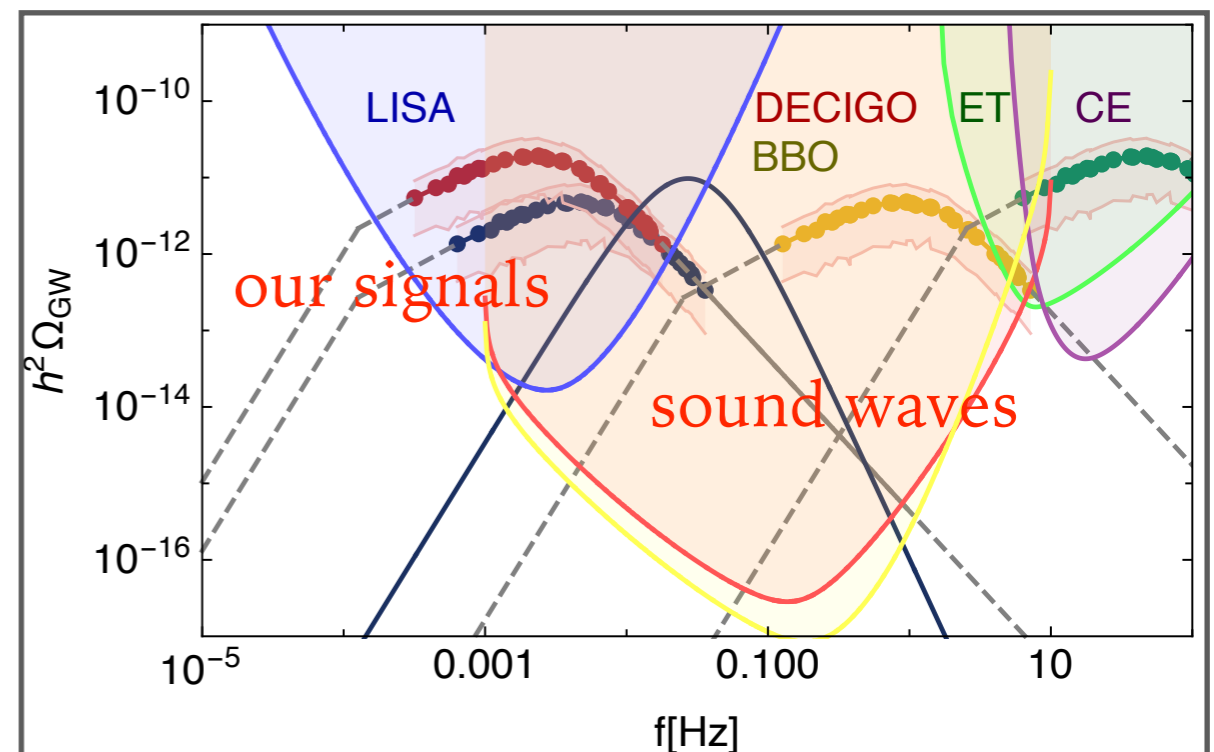
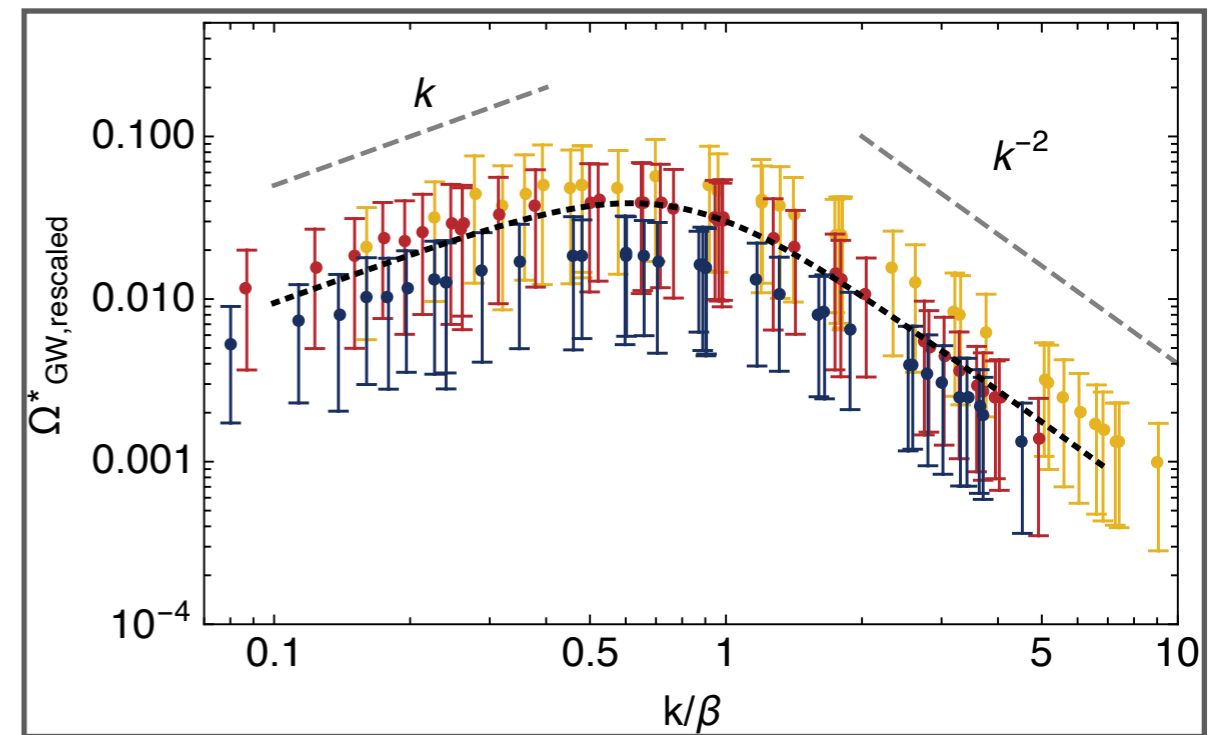
Feebly-interacting



GW SPECTRUM

► GW spectral shape is universal for different benchmark points (after normalizing by some factor)

► GW spectral shape is clearly different from sound-wave sources: it stretches over wider frequencies



DISCUSSION

► What is the essential difference?

- Sound waves: it is fluid velocity \vec{v} that superposes linearly

$$\vec{v} = \sum_{I: \text{bubbles}} \vec{v}^{(I)} \longrightarrow T_{ij} \sim wv_i v_j \neq \sum_{I: \text{bubbles}} T_{ij}^{(I)}$$

- Free-streaming particles: it is T_{ij} that superposes linearly

$$T_{ij} = \sum_{I: \text{bubbles}} T_{ij}^{(I)}$$



Introduction

*GWs from
sound waves
in FOPT*

*GWs from
FIPs
in FOPT*

Summary

SUMMARY

- We point out the missing possibility for GW sources in FOPT:

free-streaming particles

- We propose a novel GW calculation scheme ("sprinkler picture") that makes use of the linearity of GW production in each sprinkler
- Resulting GW signal is clearly different from the GW signal from the well-known sound waves