

Gravitational waves from Gauss-Bonnet-corrected single-field inflation

[based on 2308.13272 and 2108.01340 (with Shinsuke Kawai)]

Jinsu Kim

Tongji University, Shanghai, China

November 6th, 2023

at GWBSM 2023, Osaka Metropolitan University, Osaka, Japan

Inflation with Gauss-Bonnet coupling

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \frac{1}{2} g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - V(\varphi) - \frac{1}{16} \xi(\varphi) R_{\text{GB}}^2 \right]$$

$$R_{\text{GB}}^2 \equiv R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}$$

- Higher-curvature correction [Weinberg (2008)]
- 1-loop correction to some superstring models [Antoniadis, Rzos, and Tamvakis (1994)], [Rzos and Tamvakis (1994)]
- Widely studied in the context of cosmology, including slow-roll inflation [Kawai, Sakagami, and Soda (1998)], [Kawai, Sakagami, and Soda (1999)], [Kawai and Soda (1999)], [Satoh and Soda (2008)], [Satoh (2010)], [Kawai and JK (2019)], ...

$$0 = \ddot{\varphi} + 3H\dot{\varphi} + V_{,\varphi} + \frac{3}{2}H^2 (\dot{H} + H^2) \xi_{,\varphi}$$

- **Potential term \gg GB term:**
 - usual slow-roll inflation with GB coupling as a small correction
- **Potential term \approx GB term:**
 - Balance; cancellation when $V_{,\varphi}\xi_{,\varphi} < 0$
 - first part of the talk
- **Potential term \ll GB term:**
 - GB dominance
 - second part of the talk

Potential term \approx GB term with $V_{,\varphi}\xi_{,\varphi} < 0$

$$0 = \ddot{\varphi} + 3H\dot{\varphi} + V_{,\varphi} + \frac{3}{2}H^2 (\dot{H} + H^2) \xi_{,\varphi}$$
$$\implies \ddot{\varphi} + 3H\dot{\varphi} \approx 0.$$

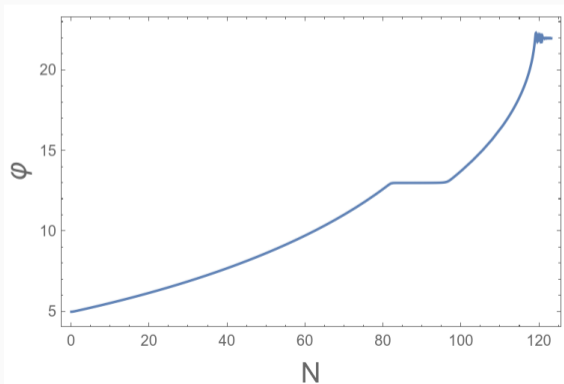
- Ultra-slow-roll inflation

[Inoue and Yokoyama (2002)], [Kinney (2005)],... see also [Kawaguchi and Tsujikawa (2022)]

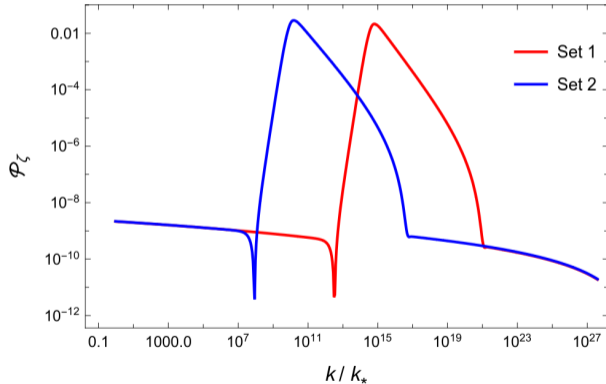
Ultra-slow-roll with Gauss-Bonnet coupling

For concreteness, let us consider

$$V = \Lambda^4 \left(1 + \cos \frac{\varphi}{f} \right), \quad \xi = \xi_0 \tanh [\xi_1 (\varphi - \varphi_c)].$$



$$\begin{aligned} & \{f[M_{\text{P}}], \varphi_c[M_{\text{P}}], \xi_0, \xi_1[M_{\text{P}}^{-1}]\} \\ & = \{7, 13.0, 6.044 \times 10^7, 15.0\} \end{aligned}$$



Curvature perturbation enhanced

- $A_s \sim 1/\epsilon_1$ or $\sim 1/\dot{\varphi}$
- production of primordial black holes
- scalar-induced gravitational waves

3 free parameters for a given f :

- φ_c : peak position
- ξ_1 : width of the peak
- ξ_0 : magnitude of the peak

Primordial black holes

When very large density fluctuations re-enter the horizon, primordial black holes may form due to the gravitational collapse. [Zel'dovich and Novikov (1967)], [Hawking (1971)], [Carr and Hawking (1974)], [Polnarev and Khlopov (1985)], ...

Abundance of the PBHs:

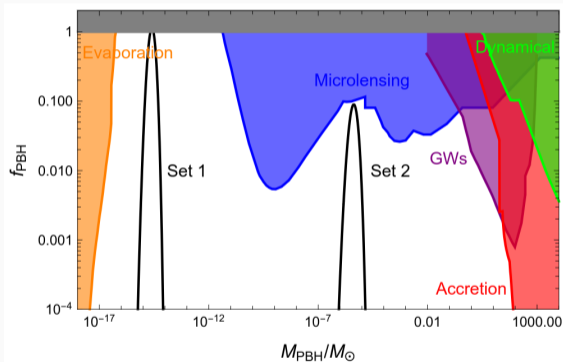
$$f_{\text{PBH}} \equiv \frac{\Omega_{\text{PBH},0}}{\Omega_{\text{DM},0}} \approx \left(\frac{\beta(M)}{3.27 \times 10^{-8}} \right) \left(\frac{0.2}{\gamma} \right)^{3/2} \left(\frac{106.75}{g_{*,f}} \right)^{1/4} \left(\frac{0.12}{\Omega_{\text{DM},0} h^2} \right) \left(\frac{M}{M_{\odot}} \right)^{-1/2}$$
$$\beta = \int_{\delta_c} d\delta \frac{1}{\sqrt{2\pi\sigma^2}} e^{-\frac{\delta^2}{2\sigma^2}}, \quad \sigma^2 = \frac{16}{81} \int_0^{\infty} \frac{dq}{q} \left(\frac{q}{k} \right)^4 W^2 \left(\frac{q}{k} \right) \mathcal{P}_{\zeta}(q)$$

Assumptions:

- Density fluctuation follows a Gaussian distribution.
- PBH formation occurs during radiation-dominated era.

Primordial black holes

constraints data : [Green and Kavanagh (2021)]



- Set 1 : $f_{\text{PBH}}^{\text{tot}} \approx 1$
- Set 2 : $f_{\text{PBH}}^{\text{tot}} \approx 0.087$

[Matarrese, Mollerach, and Bruni (1998)], [Mollerach, Harari, and Matarrese (2004)], [Ananda, Clarkson, and Wands (2007)], [Baumann, Steinhardt, Takahashi, and Ichiki (2007)], [Kohri and Terada (2018)], [Domènech (2020)]

Enhanced curvature perturbation : a source, $S_{\mathbf{k}}$, for the tensor perturbation

$$h_{\mathbf{k}}'' + 2\mathcal{H}h_{\mathbf{k}}' + k^2 h_{\mathbf{k}} = S_{\mathbf{k}}$$

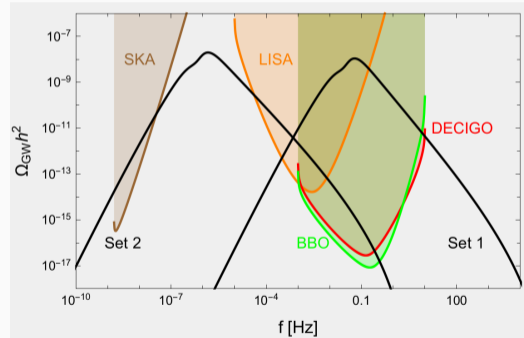
Scalar-induced second-order gravitational waves

Scalar-induced gravitational waves

[Kohri and Terada (2018)]

$$\begin{aligned}\Omega_{\text{GW},f}(k) &= \frac{1}{12} \int_0^\infty dv \int_{|1-v|}^{1+v} du \\ &\times \left(\frac{4v^2 - (1 + v^2 - u^2)^2}{4uv} \right)^2 \\ &\times \mathcal{P}_\zeta(kv) \mathcal{P}_\zeta(ku) \left(\frac{3(u^2 + v^2 - 3)}{4u^3v^3} \right)^2 \\ &\times \left[\left(-4uv + (u^2 + v^2 - 3) \log \left| \frac{3 - (u+v)^2}{3 - (u-v)^2} \right| \right)^2 \right. \\ &\left. + \pi^2 (u^2 + v^2 - 3)^2 \theta(v + u - \sqrt{3}) \right].\end{aligned}$$

sensitivity curves : [Schmitz (2021)]



$$0 = \ddot{\varphi} + 3H\dot{\varphi} + V_{,\varphi} + \frac{3}{2}H^2 (\dot{H} + H^2) \xi_{,\varphi}$$

- **Potential term \approx GB term:**
 - Balance; cancellation when $V_{,\varphi}\xi_{,\varphi} < 0$
 - SR \rightarrow USR \rightarrow SR
 - primordial black hole formation
 - scalar-induced gravitational waves
- **Potential term \ll GB term:**
 - GB dominance
 - SR \rightarrow GB domination \rightarrow SR
 - with $V_{,\varphi}\xi_{,\varphi} > 0$

$$0 = \ddot{\varphi} + 3H\dot{\varphi} + V_{,\varphi} + \frac{3}{2}H^2 (\dot{H} + H^2) \xi_{,\varphi}$$

We again consider

$$\xi = \xi_0 \tanh [\xi_1(\varphi - \varphi_c)] .$$

As $\xi_{,\varphi} \sim \text{sech}^2[\xi_1(\varphi - \varphi_c)]$, the GB term may become dominant.

φ away from φ_c

SR

φ close to φ_c

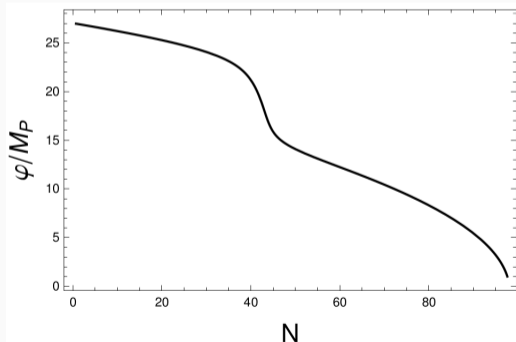
GB domination

φ away from φ_c

SR

Gauss-Bonnet domination regime

For completeness, let us consider $V = m^2\varphi^2/2$.



$$\{\xi_0, \xi_1, \varphi_c\} = \{0.196M_{\text{P}}^2/m^2, 0.5/M_{\text{P}}, 18.5M_{\text{P}}\}$$

- sudden acceleration \rightarrow deceleration
- opposite to the USR case
- surge of gravitational wave

$$v_{\mathbf{k}}'' + \left(k^2 C_t^2 - \frac{A_t''}{A_t} \right) v_{\mathbf{k}} = 0$$

$$A_t^2 \equiv a^2 \left(1 - \frac{\sigma_1}{2} \right), \quad C_t^2 \equiv 1 + \frac{a^2 \sigma_1}{2A_t^2} (1 - \sigma_2 - \epsilon_1), \quad \sigma_1 \equiv \frac{H\dot{\xi}}{M_{\text{P}}^2}, \quad \sigma_2 \equiv \frac{\dot{\sigma}_1}{H\sigma_1}$$

- Away from φ_c :

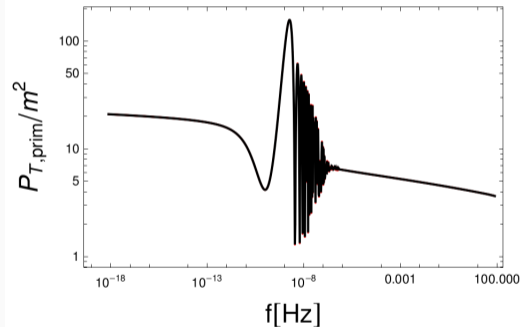
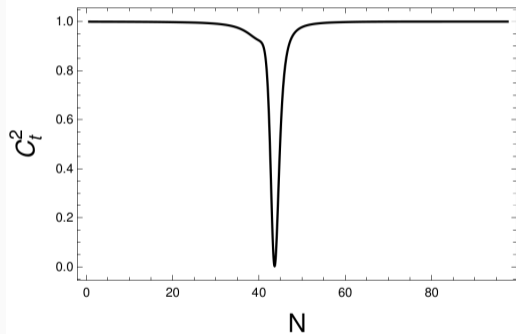
$$A_t \approx a, \quad C_t^2 \approx 1 \quad \implies \quad \text{standard tensor perturbation}$$

- Near φ_c :

$$C_t^2 \ll 1$$

Primordial tensor power spectrum

$$\{\xi_0, \xi_1, \varphi_c\} = \{0.196 M_{\text{P}}^2/m^2, 0.5/M_{\text{P}}, 18.5 M_{\text{P}}\}$$

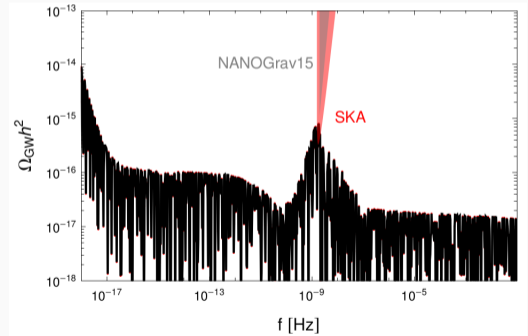


see also [Sato, Kanno, and Soda (2008)], [Guo and Schwarz (2009)], [Cai, Wang, and Piao (2015, 2016)], [Cai and Piao (2022)]

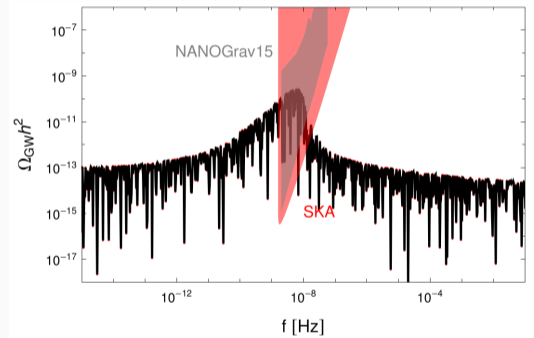
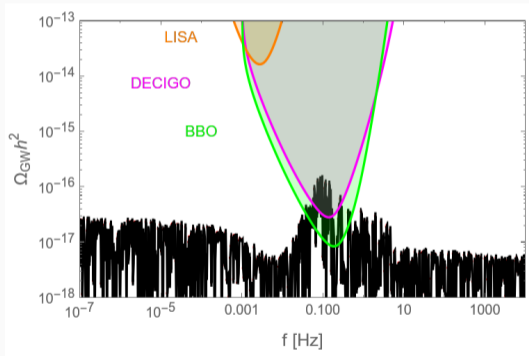
Gravitational wave spectrum

[Guzzetti, Bartolo, Liguori, and Matarrese (2016)], [Kuroyanagi, Takahashi, and Yokoyama (2015, 2021)], [Boyle and Steinhardt (2008)], ...

$$\Omega_{\text{GW}}(k) = \frac{1}{12} \left(\frac{k}{a_0 H_0} \right)^2 T^2(k) \mathcal{P}_{\text{T,prim}}$$



Gravitational wave spectrum



$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{P}}^2}{2} R - \frac{1}{2} g^{\mu\nu} \partial_\mu \varphi \partial_\nu \varphi - V(\varphi) - \frac{1}{16} \xi(\varphi) R_{\text{GB}}^2 \right]$$

Potential term \approx GB term

- Balance; $V_{,\varphi} \xi_{,\varphi} < 0$
- SR \rightarrow USR \rightarrow SR
- PBH formation
- induced gravitational waves

Potential term \ll GB term

- GB dominance; $V_{,\varphi} \xi_{,\varphi} > 0$
- SR \rightarrow GB domination \rightarrow SR
- primordial inflationary gravitational waves enhanced

Thank you