

# 超弦理論と宇宙定数問題

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- Asymmetric Orbifolds as Non-SUSY String Vacua with Vanishing CC
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# Introduction : Cosmological Constant in Superstring

# What's cosmological constant?

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Cosmological const. = vacuum energy density

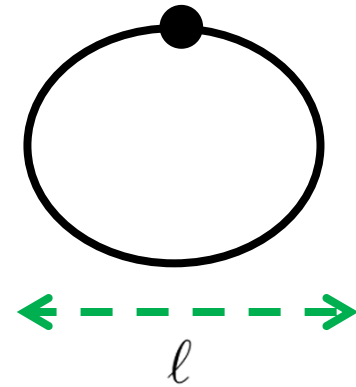
$$\Lambda_{1\text{-loop}} \sim \sum [\text{1-loop vacuum diagram}]$$

no interaction



# What's cosmological constant?

Point particle



$$\Lambda = \frac{1}{V_d} \left[ \sum_{i \in \mathcal{H}_B} - \sum_{i \in \mathcal{H}_F} \right] Z_{S^1}(m_i^2)$$
$$\equiv \left[ \sum_{i \in \mathcal{H}_B} - \sum_{i \in \mathcal{H}_F} \right] \int_{\epsilon}^{\infty} \frac{d\ell}{\ell} \int \frac{d^d p_i}{(2\pi)^d} e^{-\ell \frac{\alpha'}{2} [p_i^2 + m_i^2]}$$

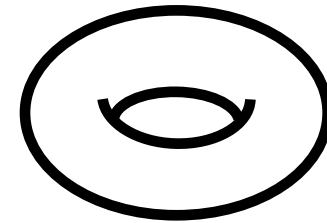
UV-cut off

Schwinger parameter  
(moduli)

# What's cosmological constant?

String theory

Torus partition function



Torus, moduli  $\mathcal{T}$

$$\begin{aligned}\Lambda &= \frac{1}{V_d} \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} \underline{Z_{\text{torus}}(\tau)} \\ &= \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} \frac{1}{\tau_2^{\frac{d-2}{2}}} \sum_{(i, \tilde{i}) \in \mathcal{H}_{\perp}} D(h_i, \tilde{h}_{\tilde{i}}) q^{h_i - \frac{1}{2}} \overline{q^{\tilde{h}_{\tilde{i}} - \frac{1}{2}}} \\ &\quad (q \equiv e^{2\pi i\tau})\end{aligned}$$

# What's cosmological constant?

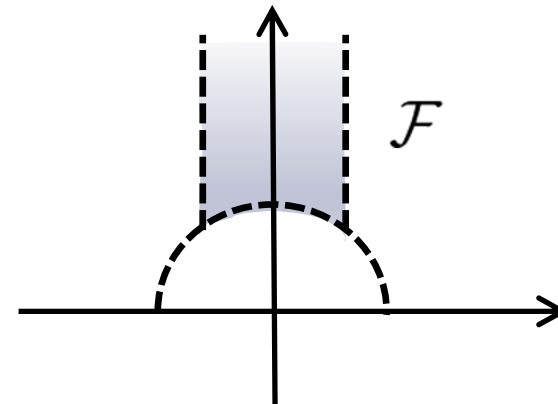
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Modular invariance

$$Z_{\text{torus}}\left(-\frac{1}{\tau}\right) = Z_{\text{torus}}(\tau), \quad Z_{\text{torus}}(\tau + 1) = Z_{\text{torus}}(\tau)$$



integration region :



# What's cosmological constant?

✂ CC for string is formally identified as that of a particle theory with an **infinite** number of mass spectrum:

$$\left( \tau = \frac{\theta + i\ell}{2\pi}, \quad \frac{\alpha'}{4} m_i^2 = h_i - \frac{1}{2} = \tilde{h}_{\tilde{i}} - \frac{1}{2} \right)$$
$$\Lambda = \int_0^\infty \frac{d\ell}{\ell} \int_{-\pi}^\pi \frac{d\theta}{2\pi} (2\pi\alpha'\ell)^{-\frac{d}{2}}$$
$$\times \sum_{i, \tilde{i}} D(h_i, \tilde{h}_{\tilde{i}}) \exp \left[ - \left( h_i + \tilde{h}_{\tilde{i}} - 1 \right) \ell + i\theta \left( h_i - \tilde{h}_{\tilde{i}} \right) \right]$$
$$\propto \int_{\mathcal{S}} \frac{d^2\tau}{\tau_2^2} \frac{1}{\tau_2^{\frac{d-2}{2}}} \sum_{i, \tilde{i}} D(h_i, \tilde{h}_{\tilde{i}}) q^{h_i - \frac{1}{2}} \overline{q^{\tilde{h}_{\tilde{i}} - \frac{1}{2}}},$$

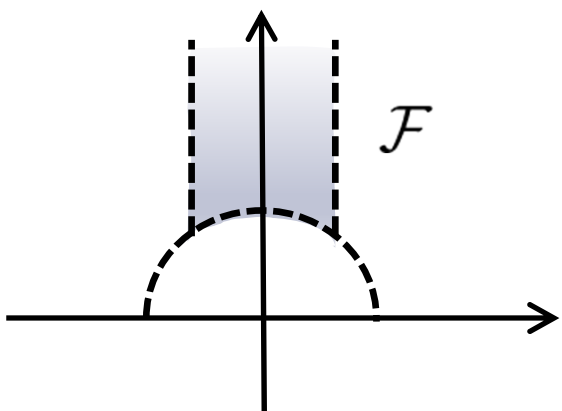
But, the integration region is different.



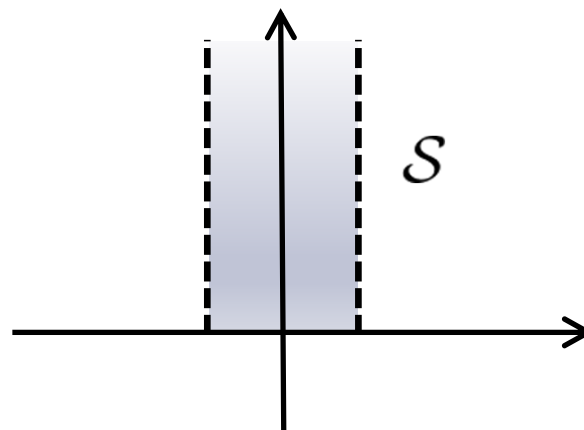
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Integration region of moduli  
(Schwinger parameter)

String



Particle



# 'cosmological constant problem'

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$$\Lambda_{\text{observation}} \ll M_{\text{SUSY breaking}}^4$$



suggests a vacuum with (nearly)  
vanishing cosmological constant  
without SUSY?

## Non-SUSY with $\Lambda = 0$ ?

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Two possibilities

(1)  $Z(\tau) \equiv 0$  , but no supercharges exist.

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(2)  $Z(\tau) \neq 0$  , but  $\Lambda = 0$

**Asymmetric Orbifolds**  
as  
**Non-SUSY String Vacua with  
Vanishing Cosmological Constant**

# How to break SUSY?

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Orbifolding by  $(-1)^{F_L}$

$(F_L : \text{left-moving space-time fermion number})$



Removes massless Ramond states in the untwisted sector.  
**(break SUSY?)**



But, new massless Ramond states emerge in the twisted sector. **(eventually, SUSY)**

# How to break SUSY?

Orbifolding by  $(-1)^{F_L} \otimes \mathcal{T}_{2\pi R}$

‘Scherk-Schwarz  
Compactification’

(see also Prof. Itoyama’s talk)

shift  
operator



Lightest Ramond states in the twisted sectors  
get massive. (No supercharge can appear.)



SUSY breaking

# Non-SUSY Model with $\Lambda = 0$

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Non-SUSY string vacua (type II)  
with vanishing  
1-loop cosmological constant

[Kachru, Kumar, Silverstein 1998], [Kachru, Silverstein 1998],  
[Kachru, Silverstein 1998]

Asymmetric orbifold defined by the orbifold group:

$$G = \langle f, g \rangle$$

$f, g$  do not commute.



Non-abelian  
orbifold

$f$  → includes  $(-1)^{F_R}$  and (chiral) shift op.

→ break all the right-moving SUSY

$g$  → includes  $(-1)^{F_L}$  and (chiral) shift op.

→ break all the left-moving SUSY



**Totally Non-SUSY**



Nevertheless, the 1-loop cosmological constant  
(torus partition function) strictly vanishes.

Bose-fermi cancelation  
without SUSY

Crucial point :

$$\alpha \begin{array}{|c|} \hline \square \\ \hline \end{array} \beta \quad (\alpha, \beta \in G)$$

is well-defined, only when  $\alpha$  commutes with  $\beta$



Clever, but looks complicated...

## Closely related studies :

[Harvey 1998]

[Shiu, Tye 1998]

[Blumenhagen, Gorlich 1998]

[Angelantonj, Antoniadis, Forger 1999]

[Aoki, D'Hoker, Phong 2003 ]

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# Non-SUSY Model with $\Lambda = 0$

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## My studies

[Sato, Y.S, Wada 2015]

[Y.S, Wada 2016]

[Sato, Y.S, Uetoko 2017]

[Y.S, Uetoko 2018]

[Aoyama, Y.S 2020]

[Aoyama, Y.S to appear]

# Non-SUSY Model with $\Lambda = 0$

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[Satoh, Y.S, Wada 2015]

- Realized by simpler (asymmetric) orbifold actions.  
(only single generator in the orbifold group)
- Each building block  $Z_{(a,b)}(\tau)$  separately vanishes.  
(bose-fermi cancellation in each sector)
- No supercharges (globally defined on the total Hilbert space) exist.

# Non-SUSY Model with $\Lambda = 0$

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$$Z(\tau) = \sum_{a,b} \underline{Z_{(a,b)}(\tau)} \equiv \sum_{a,b} b \square_a$$

$$Z_{(a,b)}\left(-\frac{1}{\tau}\right) = Z_{(b,-a)}(\tau)$$

$$Z_{(a,b)}(\tau + 1) = Z_{(a,a+b)}(\tau)$$

← ‘modular covariance’

# Non-SUSY Model with $\Lambda = 0$

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[Aoyama, Y.S 2020]

[Aoyama, Y.S to appear]

- Asymmetric orbifolds of **Gepner models**
- Each building block  $Z_{(a,b)}(\tau)$  does **not** necessarily vanish.

**(bose-fermi cancellation among different twisted sectors)**

# Non-SUSY Vacua with Vanishing Cosmological Constant : part 2

[Satoh, Y.S work in progress]

# Non-SUSY Model with $\Lambda = 0$ : part 2

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Two possibilities

(1)  $Z(\tau) \equiv 0$  , but no supercharges exist.

(2)  $Z(\tau) \neq 0$  , but  $\Lambda = 0$

no bose-fermi cancellation



# Non-SUSY Model with $\Lambda = 0$ : part 2

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[Moore 1987] 'Atkin-Lehner Symmetry'

$$\Lambda = \langle \psi_1 | \psi_2 \rangle$$

If  $\psi_1, \psi_2$  have **opposite** 'parity',

$\Lambda$  should vanish.



Beautiful idea, but difficulty in concrete model building ...

# Non-SUSY Model with $\Lambda = 0$ : part 2

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Our idea :

$$Z_{\text{total}}(\tau) = \sum_{w,m \in \mathbb{Z}} Z_{(w,m)}(\tau)$$



**Modular invariant**

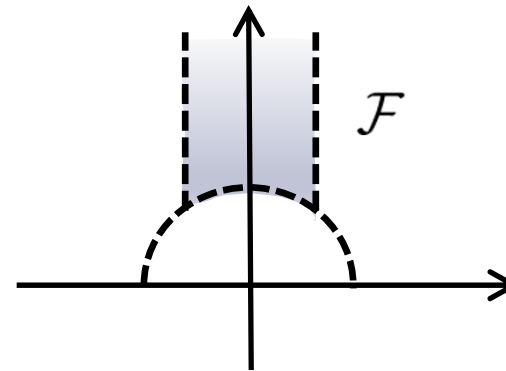
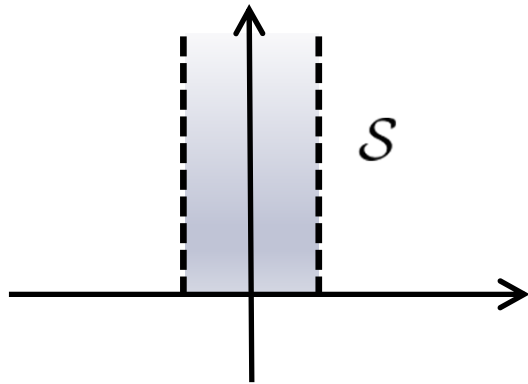
$(w, m)$  : 'winding numbers'  
(  $SL(2, \mathbb{Z})$  -doublet )

# 'Polchinski's trick'

[Polchinski 1986]


$$\int_{\mathcal{S}} \frac{d^2\tau}{\tau_2^2} \sum_{m \in \mathbb{Z}} Z_{(0,m)}(\tau) = \int_{\mathcal{F}} \frac{d^2\tau}{\tau_2^2} \sum_{w,m \in \mathbb{Z}} Z_{(w,m)}(\tau)$$

( assume  $Z_{(0,0)}(\tau) \equiv 0$  )



Assume the situation :

$$\sum_{m \in \mathbb{Z}} Z_{(0,m)}(\tau) \neq 0, \text{ but } \sum_{w,m \in \mathbb{Z}} Z_{(w,m)}(\tau) \equiv 0$$

  $'Z_{\text{particle}}(\tau)' \left( \equiv \sum_{m \in \mathbb{Z}} Z_{(0,m)}(\tau) \right) \neq 0$

but  $\Lambda \equiv \int_{\mathcal{S}} \frac{d^2\tau}{\tau_2^2} 'Z_{\text{particle}}(\tau)' = 0$

Correct integration region  
for the particle theory !



interpretable as a **particle theory** with  $\Lambda = 0$  ,  
**without the bose-fermi cancellation** ?  
(including infinite number of mass spectrum)

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Starting with various SUSY string vacua,  
we can realize this situation by implementing  
**asymmetric orbifolding**

$$Z_{(0,m)}(\tau) \sim \text{'untwisted sector'}$$

$$Z_{(w,m)}(\tau) \ (w \neq 0) \sim \text{'twisted sector'}$$

# Summary

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- Non-SUSY string vacua with  $Z(\tau) \equiv 0$ 
  - ➔ can be constructed by asymmetric (‘non-geometric’) orbifold
- Vanishing CC with  $Z(\tau) \neq 0$ 
  - ➔ very difficult in string theory,  
can be realized as particle vacua ?

ご清聴ありがとうございました。