

# Regge conformal blocks from the Rindler-AdS black hole and the pole-skipping phenomena

**Mitsuhiro Nishida**

(Gwangju Institute of Science and Technology, Korea)

[[arXiv:2105.07778](https://arxiv.org/abs/2105.07778)]

with Keun-Young Kim, Kyung-Sun Lee

# Regge limit is constrained by consistencies of theories

Regge limit in scattering amplitude

Regge limit of  
bulk S-matrix

$$M(s, t) \quad (s \rightarrow \infty, \text{fixed } t)$$

Regge limit in CFTs

Regge limit of  
Mellin amplitude

$$\mathcal{M}(s, t) \quad (s \rightarrow \infty, t \text{ fixed})$$

Regge limit  
of CFT correlators

$$\mathcal{A}(z, \bar{z}) \quad (1 - z) \rightarrow e^{-2\pi i} (1 - z), \bar{z} \text{ fixed}$$
$$z, \bar{z} \rightarrow 0, \frac{\bar{z}}{z} \text{ fixed}$$

# Regge limit in CFTs is related to out-of-time-order correlator (OTOC)

OTOC is a measure of quantum chaos.

$$\frac{\langle W(t, \mathbf{d}) V(0, 0) W(t, \mathbf{d}) V(0, 0) \rangle}{\langle W(t, \mathbf{d}) W(t, \mathbf{d}) \rangle \langle V(0, 0) V(0, 0) \rangle} \sim 1 - \varepsilon e^{\lambda_L (t - \mathbf{d}/v_B)} + \dots$$

[A. Larkin, Y. N. Ovchinnikov, 1969], [A. Kitaev, 2014]

Regge limit corresponds to an analytic continuation from Euclidean correlator to OTOC

[D. A. Roberts, D. Stanford, 2014] [E. Perlmutter, 2016]

Lyapunov exponent  $\lambda_L$  is bounded:  $\lambda_L \leq 2\pi/\beta$ , which is related to Regge growth bound of S-matrix

[J. Maldacena, S. H. Shenker, D. Stanford, 2015],

[D. Chandorkar, S. D. Chowdhury, S. Kundu, S. Minwalla, 2021]

# Regge conformal block with $T_{\mu\nu}$ exchange captures $\lambda_L$ and $\nu_B$ of holographic CFTs

Conformal block expansion in  $d$ -dim CFTs

$$\langle W(x_1)W(x_2)V(x_3)V(x_4) \rangle = \frac{1}{x_{12}^{2\Delta_W} x_{34}^{2\Delta_V}} \sum_{\mathcal{O}} C_{WW\mathcal{O}} C_{VV\mathcal{O}} G_{\Delta,\ell}(z, \bar{z})$$

Regge conformal blocks in OTOCs

$$G_{\Delta,\ell}^{\text{Regge}}(z, \bar{z}) \propto \underline{e^{(\ell-1)t - (\Delta-1)d}} {}_2F_1 \left( \Delta - 1, \frac{d}{2} - 1, \Delta + 1 - \frac{d}{2}; e^{-2d} \right)$$

$e^{t - (d-1)d}$  with  $T_{\mu\nu}$  exchange ( $\ell = 2, \Delta = d$ )

captures  $\lambda_L = \frac{2\pi}{\beta}, \nu_B = \frac{1}{d-1}$  in the Rindler spacetime  
( $\beta = 2\pi$ )

[D. A. Roberts, D. Stanford, 2014] [E. Perlmutter, 2016]

# Proposal of pole-skipping phenomena

$\lambda_L$  and  $v_B$  in theories with Einstein gravity duals are related to a pole-skipping point  $(\omega_*, k_*)$  in  $G_{T_{00}T_{00}}(\omega, k)$ .

$$e^{-i\omega_* t + ik_* \mathbf{d}} = e^{\lambda_L (t - \mathbf{d}/v_B)}$$

[S. Grozdanov, K. Schalm, V. Scopelliti, 2017], [M. Blake, H. Lee, H. Liu, 2018],

**Pole-skipping points**  $(\omega_*, k_*)$  of  $G(\omega, k)$

Intersections between lines of poles and lines of zeros  
in momentum space of Green's function

e.g.  $G(\omega, k) = \frac{\omega + k}{\omega - k}$

Pole-skipping point

$$\omega_* = k_* = 0$$

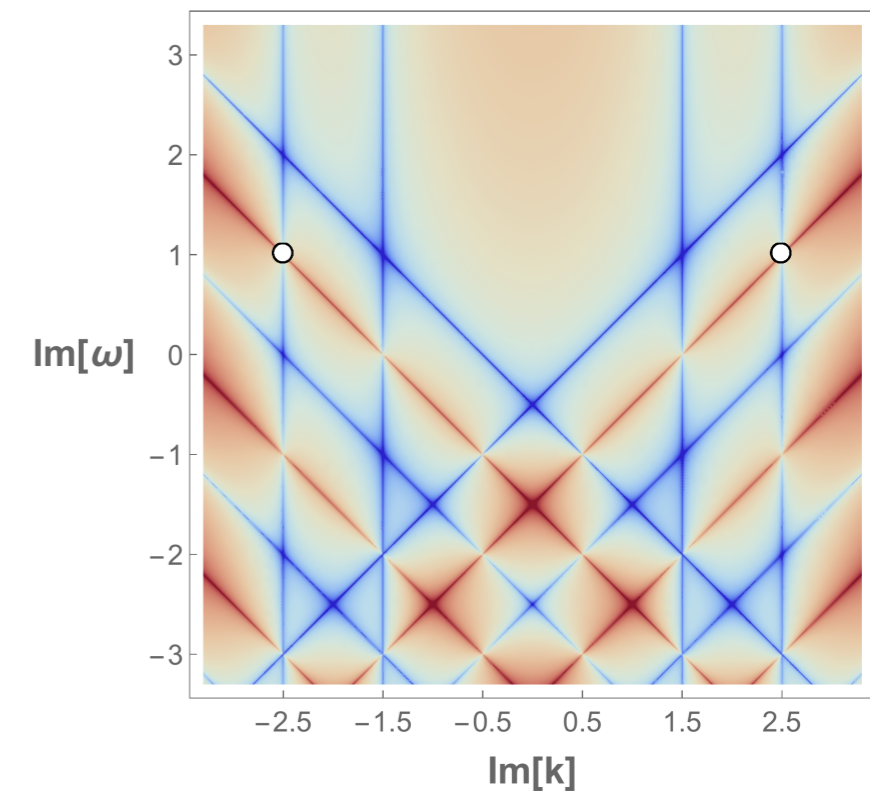
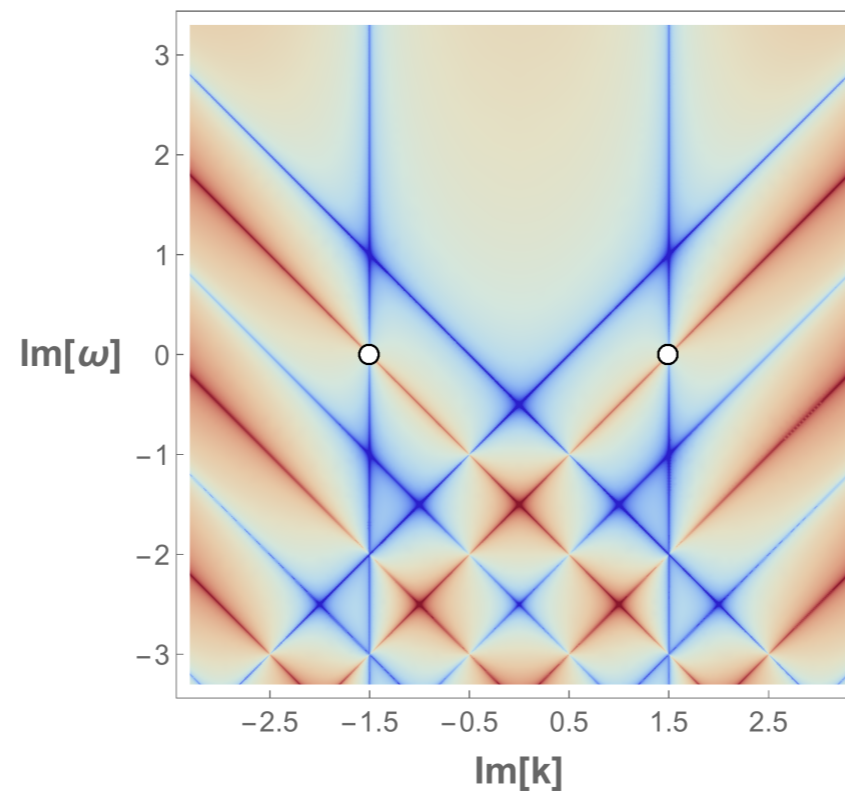
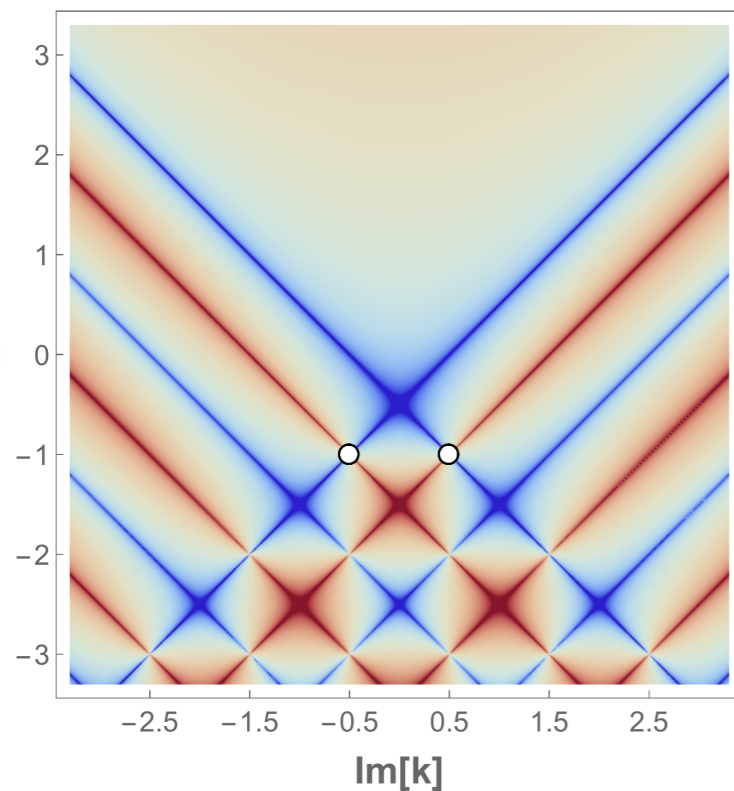
# Leading pole-skipping points are related to Regge conformal blocks

$$G_{\Delta, \ell}^{\text{Regge}} \propto e^{(\ell-1)t - (\Delta-1)d}$$

$\ell = 0$

$\ell = 1$

$\ell = 2$



[Y. Ahn, V. Jahnke, H.-S. Jeong,  
K.-Y. Kim, K.-S. Lee, MN, 2020]

[F. M. Haehl, W. Reeves,  
M. Rozali, 2019]

red lines: poles    blue lines: zeros

intersections: pole-skipping points

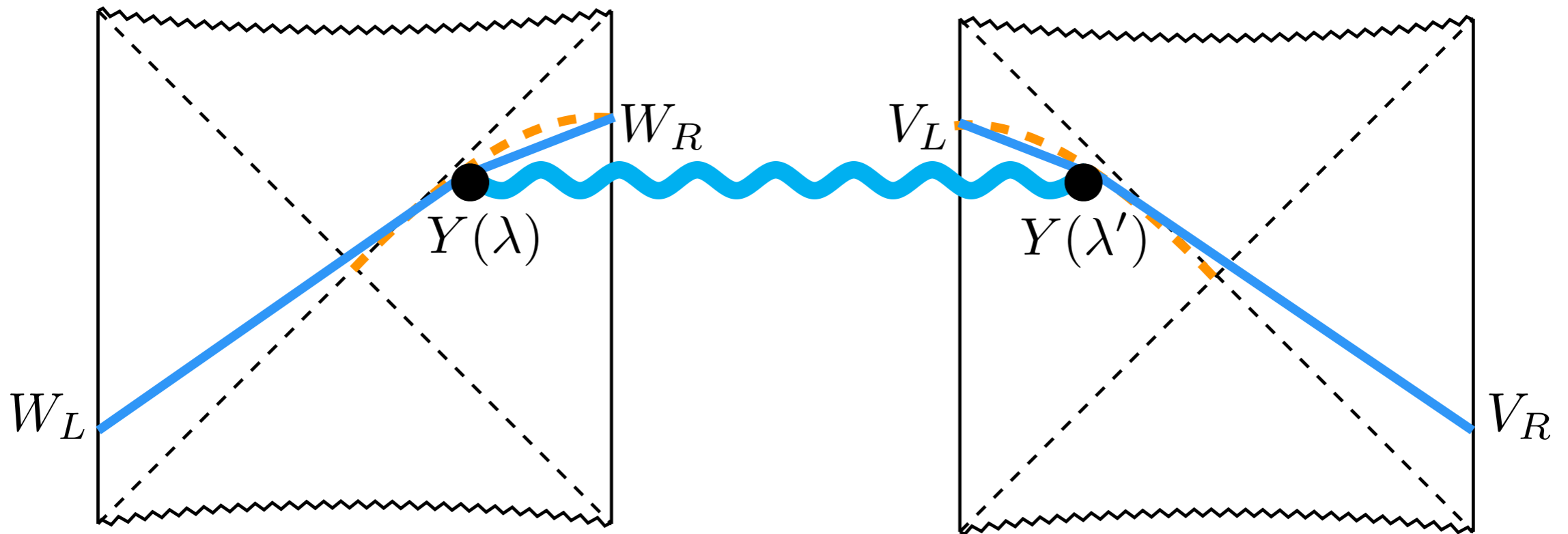
# Our motivation

Holographic understanding why pole-skipping points capture Regge conformal blocks

- Near-horizon analysis is a holographic method to detect pole-skipping points [M. Blake, R. A. Davison, S. Grozdanov, H. Liu, 2018]
- Gravity duals of Regge conformal or OPE blocks were studied [L. Cornalba, M. S. Costa, J. Penedones, R. Schiappa, 2006]  
[N. Afkhami-Jeddi, T. Hartman, S. Kundu, A. Tajdini, 2017]  
[N. Kobayashi, T. Nishioka, Y. Okuyama, 2020]
- We want to connect Regge conformal blocks and near-horizon analysis in Rindler-AdS BH

# Our result 1

We construct diagrams in the Rindler-AdS BH that have the Regge behaviors of conformal blocks



Half-geodesic Witten diagrams

$$\begin{aligned} \mathcal{W}_{\Delta, \ell}^{\mathcal{R}} &:= \int_{\gamma_W^R} d\lambda \int_{\gamma_V^L} d\lambda' G_{b\partial} (Y(\lambda), W_L; \Delta_W) G_{b\partial} (Y(\lambda), W_R; \Delta_W) G_{bb} \left( Y(\lambda), Y(\lambda'); \frac{dY(\lambda)}{d\lambda}, \frac{dY(\lambda')}{d\lambda'}; \Delta, \ell \right) \\ &\quad \times G_{b\partial} (Y(\lambda'), V_L; \Delta_V) G_{b\partial} (Y(\lambda'), V_R; \Delta_V) \\ &\simeq \frac{(\mathcal{C}_{\Delta_W, 0} \mathcal{C}_{\Delta_V, 0})^2 \mathcal{C}_{\Delta, 0}}{2^{2(\Delta_W + \Delta_V) + \ell} (\Delta - 1)} \log \left( \frac{1}{\epsilon} \right) e^{(\ell-1)t_R - (\Delta-1)\mathbf{d}} {}_2F_1 \left( \Delta - 1, \frac{d}{2} - 1, \Delta + 1 - \frac{d}{2}; e^{-2\mathbf{d}} \right) \end{aligned}$$



# Our result 2

Near-horizon analysis in the Rindler-AdS BH captures equations of the Regge behaviors with any integer spin

$$(\nabla_{\mu}\nabla^{\mu} - \Delta(\Delta - d) + \ell)h_{v\dots v}(v, r, \mathbf{x}) = 0$$

↓ Near-horizon analysis

$$e^{(\ell-1)t} \left[ \square_{\mathbb{H}} - (\Delta - 1)(\Delta - d + 1) \right] h_{v\dots v} \Big|_{r \rightarrow 1} = 0$$

↓

↓

$$G_{\Delta, \ell}^{\text{Regge}}(z, \bar{z}) \propto \underline{e^{(\ell-1)t_R - (\Delta-1)d} {}_2F_1\left(\Delta - 1, \frac{d}{2} - 1, \Delta + 1 - \frac{d}{2}; e^{-2d}\right)}$$

Integrals of the bulk-to-bulk propagators  
are related to the near-horizon analysis

# Summary

- Pole-skipping points in momentum Green's functions are related to Regge conformal blocks in OTOCs
- We construct half-geodesic Witten diagrams on Rindler-AdS BH as gravity duals of Regge conformal blocks
- We derive equations for Regge conformal blocks with integer spin from near-horizon analysis on Rindler-AdS BH

# Future directions

- Fermion exchange
- Gravity duals with retarded or advanced propagators
- CFT's picture of the relation between Regge conform blocks and pole-skipping points
- Sub-leading analysis of the pole-skipping points and the half-geodesic Witten diagrams